Double Obstacle Limit for a Navier-Stokes/Cahn-Hilliard System

Helmut Abels

Faculty of Mathematics University of Regensburg

March 10, 2010

Diffuse Interface Models

We consider two (macroscopically) immiscible incompressible, viscous fluids like oil and water.

Classical Models: Interface is a two-dimensional surface.

Surface tension is proportional to the mean curvature.



Diffuse Interface Models

We consider two (macroscopically) immiscible incompressible, viscous fluids like oil and water.

Classical Models: Interface is a two-dimensional surface.

Surface tension is proportional to the mean curvature.



But: Sharp Interface is an idealization (van der Waals). Fluid mix in a thin interfacial region.

Phase Separation/Cahn-Hilliard Equation

We consider two partly miscible components, e.g. Al-/Ni-atoms in a melted alloy, oil and water.

Let $c_j \colon \Omega \to \mathbb{R}$ be the concentration of the component j = 1, 2, $c = c_1 - c_2$, and let

$$E_{\varepsilon}(c) = \frac{\varepsilon}{2} \int_{\Omega} |\nabla c(x)|^2 \, dx + \varepsilon^{-1} \int_{\Omega} f(c(x)) \, dx$$

be the free energy of the mixture, where $\Omega \subseteq \mathbb{R}^d$, $d=1,2,3, \ \varepsilon>0$ and

 $f: \mathbb{R} \to [0,\infty)$ with $f(c) = 0 \Leftrightarrow c = \pm 1$.



Example:
$$f(c) = \frac{1}{8}(1 - c^2)^2$$

Phase Separation/Cahn-Hilliard Equation

We consider two partly miscible components, e.g. Al-/Ni-atoms in a melted alloy, oil and water.

Let $c_j \colon \Omega \to \mathbb{R}$ be the concentration of the component j = 1, 2, $c = c_1 - c_2$, and let

$$E_{\varepsilon}(c) = \frac{\varepsilon}{2} \int_{\Omega} |\nabla c(x)|^2 \, dx + \varepsilon^{-1} \int_{\Omega} f(c(x)) \, dx$$

be the free energy of the mixture, where $\Omega\subseteq\mathbb{R}^d$, $d=1,2,3,\ arepsilon>0$ and

 $f : \mathbb{R} \to [0,\infty)$ with $f(c) = 0 \Leftrightarrow c = \pm 1$.



Example: $f(c) = \frac{1}{8}(1 - c^2)^2$

 H^{-1} -gradient flow of E_{ε} describes dynamics of phase separation:

$$\partial_t c = \Delta \mu$$
 in $\Omega \times (0, \infty)$ (1)
 $\mu = \varepsilon^{-1} f'(c) + \varepsilon \Delta c$ in $\Omega \times (0, \infty)$ (2)



Choice of Free Energy Density(I)

Typical choice: Smooth double well potential as e.g. $f(c) = \frac{1}{8}(1-c^2)^2$. Then the optimal profile of a diffuse interface is

$$c_0(x) = anh rac{x}{2arepsilon}, \qquad x \in \mathbb{R},$$



which minimizes E_{ε} in the case $\Omega = \mathbb{R}$ with constraint $c(x) \rightarrow_{x \rightarrow \pm \infty} \pm 1$. Note: $c_0(x) \in (-1, 1)$ for all $x \in \mathbb{R}$.

Choice of Free Energy Density(I)

Typical choice: Smooth double well potential as e.g. $f(c) = \frac{1}{8}(1-c^2)^2$. Then the optimal profile of a diffuse interface is

$$c_0(x) = \tanh \frac{x}{2\varepsilon}, \qquad x \in \mathbb{R},$$



which minimizes E_{ε} in the case $\Omega = \mathbb{R}$ with constraint $c(x) \rightarrow_{x \rightarrow \pm \infty} \pm 1$. Note: $c_0(x) \in (-1, 1)$ for all $x \in \mathbb{R}$.

Problem: For smooth f solutions c(x, t) of Cahn-Hilliard system

$$\partial_t c = \Delta \mu \qquad \text{in } \Omega \times (0, \infty)$$
$$\mu = \varepsilon^{-1} f'(c) + \varepsilon \Delta c \quad \text{in } \Omega \times (0, \infty)$$

might not stay in [-1, 1]!

Choice of Free Energy (II)

In the following we consider the logarithmic free energy density

$$f_{ heta}(c) = egin{cases} heta((1-c)\log(1-c)+(1+c)\log(1+c)) - heta_c c^2, & ext{if } c\in[-1,1], \ +\infty & ext{else.} \end{cases}$$

for some $0 < \theta < \theta_c$ and $\nu(c) > 0$ on [-1,1], cf. Cahn & Hilliard '58, Elliott & Luckhaus '91.

Note:





Elliott & Luckhaus '91, Debussche & Dettori '95, Kenmochi et al. '95: Existence of unique solutions of (1)-(2) such that $c(x, t) \in (-1, 1)$. Alternative proofs: Miranville & Zelik '04, A. & Wilke '07

Helmut Abels (U Regensburg)

Double Obstacle Limit

Double Obstacle/Deep Quench Limit (I)

We have

$$f_{\theta}(c) \to_{\theta \to 0} f_{0}(c) = \mathbf{I}_{[-1,1]}(c) - \frac{\theta_{c}}{2}c^{2}, \quad \mathbf{I}_{[-1,1]}(c) := \begin{cases} 0 & \text{if } c \in [-1,1], \\ +\infty & \text{else.} \end{cases}$$

The optimal profile for E_{ε} with f_0 and $\varepsilon = \theta_c = 1$ is

$$c_0(x) = \begin{cases} -1 & \text{if } c < -\frac{\pi}{2} \\ \sin x & \text{if } c \in [-\frac{\pi}{2}, \frac{\pi}{2}] \\ 1 & \text{if } c > \frac{\pi}{2} \end{cases}$$

Double Obstacle/Deep Quench Limit (II)

The optimal profile

$$c_0(x) = \begin{cases} -1 & \text{if } c < -\frac{\pi}{2} \\ \sin x & \text{if } c \in [-\frac{\pi}{2}, \frac{\pi}{2}] \\ 1 & \text{if } c > \frac{\pi}{2} \end{cases}$$

solves the differential inclusion

$$c''(x) + c(x) \in \partial I_{[-1,1]}(c(x)) = \begin{cases} (-\infty, 0] & \text{if } c(x) = -1\\ \{0\} & \text{if } c(x) \in (-1,1)\\ [0,\infty) & \text{if } c(x) = 1\\ \emptyset & \text{else} \end{cases}$$

Elliott & Luckhaus '91: Solutions of Cahn-Hilliard system (1)-(2) with f_{θ} converge as $\theta \to 0$ to solution of

$$\partial_t c = \Delta \mu \qquad \text{in } \Omega \times (0, T)$$
$$\mu + \varepsilon \Delta c + \varepsilon^{-1} \theta_c c \in \partial I_{[-1,1]}(c(x)) \qquad \text{in } \Omega \times (0, T)$$

1

Diffuse Interface Model in the Case of Matched Densities

If the densities of the fluids are the same, then one can derive:

$$\partial_{t}v + v \cdot \nabla v - \underbrace{\operatorname{div}(\nu(c)Dv)}_{\text{inner friction}} + \nabla p = \underbrace{-\varepsilon \operatorname{div}(\nabla c \otimes \nabla c)}_{\text{surface tension}} \text{ in } \Omega \times (0,\infty) \quad (3)$$

$$\operatorname{div} v = 0 \qquad \text{ in } \Omega \times (0,\infty) \quad (4)$$

$$\partial_{t}c + v \cdot \nabla c = m\Delta\mu \qquad \text{ in } \Omega \times (0,\infty) \quad (5)$$

$$\mu = -\varepsilon\Delta c + \varepsilon^{-1}f_{\theta}'(c) \text{ in } \Omega \times (0,\infty) \quad (6)$$
where $Dv = \frac{1}{2}(\nabla v + \nabla v^{T}), \ \Omega \subset \mathbb{R}^{d}$ is a bounded smooth domain,

where $Dv = \frac{1}{2}(\nabla v + \nabla v')$, $\Omega \subset \mathbb{R}^{d}$ is a bounded smooth domain, together with boundary and initial conditions. Derivation: Hohenberg & Halperin '74, Gurtin et al. '96 Analytical results: Starovoitov '93, Boyer '03, X. Feng '06, A. '07/'09

Diffuse Interface Model in the Case of Matched Densities

If the densities of the fluids are the same, then one can derive:

$$\partial_{t}v + v \cdot \nabla v - \underbrace{\operatorname{div}(\nu(c)Dv)}_{\text{inner friction}} + \nabla p = \underbrace{-\varepsilon \operatorname{div}(\nabla c \otimes \nabla c)}_{\text{surface tension}} \text{ in } \Omega \times (0,\infty) \quad (3)$$
$$\operatorname{div} v = 0 \qquad \text{ in } \Omega \times (0,\infty) \quad (4)$$
$$\partial_{t}c + v \cdot \nabla c = m\Delta\mu \qquad \text{ in } \Omega \times (0,\infty) \quad (5)$$
$$\mu = -\varepsilon\Delta c + \varepsilon^{-1}f_{\theta}'(c) \text{ in } \Omega \times (0,\infty) \quad (6)$$
where $Dv = \frac{1}{2}(\nabla v + \nabla v^{T}) \quad \Omega \subset \mathbb{R}^{d}$ is a bounded smooth domain

where $Dv = \frac{1}{2}(\nabla v + \nabla v')$, $\Omega \subset \mathbb{R}^{a}$ is a bounded smooth domain, together with boundary and initial conditions. Derivation: Hohenberg & Halperin '74, Gurtin et al. '96 Analytical results: Starovoitov '93, Boyer '03, X. Feng '06, A. '07/'09 Energy dissipation: For sufficiently smooth solutions we have

$$\frac{d}{dt}E(c(t), v(t)) = -\int_{\Omega} \nu(c)|Dv|^{2} dx - \int_{\Omega} m|\nabla\mu|^{2} dx \quad \text{with}$$

$$E(c(t), v(t)) = \frac{\varepsilon}{2} \int_{\Omega} |\nabla c(x)|^{2} dx + \varepsilon^{-1} \int_{\Omega} f_{\theta}(c(x)) dx + \int_{\Omega} \frac{|v(t)|^{2}}{2} dx$$

Theorem (Existence, Regularity, Uniqueness, A. '07/'09)

Let $d = 2, 3, \theta > 0$. For every $v_0 \in L^2_{\sigma}(\Omega)$, $c_0 \in H^1(\Omega)$ with $E(c_0, v_0) < \infty$ there is a weak solution (v, c, μ) of (3)-(6), which satisfies

$$(v, \nabla c) \in L^{\infty}(0, \infty; L^{2}(\Omega)), \quad (\nabla v, \nabla \mu) \in L^{2}(0, \infty; L^{2}(\Omega)),$$

 $\nabla^{2}c, f_{\theta}'(c) \in L^{2}_{loc}([0, \infty); L^{6}(\Omega)).$

In particular, $c(t,x) \in (-1,1)$ a.e. Moreover, $c \in BUC([0,\infty); W^1_q(\Omega))$ with q > d. For (v_0, c_0) sufficiently smooth:

- If d = 2, then the weak solution is unique and regular.
- ② If d = 3, there are some $0 < T_0 < T_1 < \infty$ such that the weak solution is regular and (locally) unique on $(0, T_0)$ and $[T_1, \infty)$.

Theorem (Existence, Regularity, Uniqueness, A. '07/'09)

Let $d = 2, 3, \theta > 0$. For every $v_0 \in L^2_{\sigma}(\Omega)$, $c_0 \in H^1(\Omega)$ with $E(c_0, v_0) < \infty$ there is a weak solution (v, c, μ) of (3)-(6), which satisfies

$$(v, \nabla c) \in L^{\infty}(0, \infty; L^{2}(\Omega)), \quad (\nabla v, \nabla \mu) \in L^{2}(0, \infty; L^{2}(\Omega)),$$

 $\nabla^{2}c, f'_{\theta}(c) \in L^{2}_{loc}([0, \infty); L^{6}(\Omega)).$

In particular, $c(t,x) \in (-1,1)$ a.e. Moreover, $c \in BUC([0,\infty); W_q^1(\Omega))$ with q > d. For (v_0, c_0) sufficiently smooth:

- If d = 2, then the weak solution is unique and regular.
- ② If d = 3, there are some $0 < T_0 < T_1 < \infty$ such that the weak solution is regular and (locally) unique on $(0, T_0)$ and $[T_1, \infty)$.

Theorem (Uniform Bounds, A. '09)

The solutions $(v, c, \mu) = (v^{\theta}, c^{\theta}, \mu^{\theta})$, $\theta \in (0, 1)$ are uniformly bounded in the function spaces above.

Structure of the Proof

First study the separate systems:

- Cahn-Hilliard equation with convection and singular potential (based on $E_{\varepsilon}(c) = E_{0,\theta}(c) \frac{\theta_c}{2} \|c\|_2^2$ with $E_{0,\theta}$ convex)
- (Navier-)Stokes system with variable viscosity

Structure of the Proof

First study the separate systems:

- Cahn-Hilliard equation with convection and singular potential (based on $E_{\varepsilon}(c) = E_{0,\theta}(c) \frac{\theta_c}{2} \|c\|_2^2$ with $E_{0,\theta}$ convex)
- (Navier-)Stokes system with variable viscosity

Existence of weak solutions:

Approximation and compactness argument

Higher Regularity: Use regularity results for separate systems

Uniqueness: Gronwall's inequality once $c \in L^{\infty}(0, T; C^{1}(\overline{\Omega}))$ and $v \in L^{\infty}(0, T; W^{1}_{s}(\Omega)), s > d$.

Crucial ingredient for higher regularity: A priori estimate for $c \in BUC([0,\infty); W^1_q(\Omega)), q > d!$

A priori Estimates for c

 W_r^2 -estimate for c: Formally multiply

$$\begin{split} \mu(x,t) &= -\Delta c(x,t) + f'_{\theta}(c(x,t)) \\ f'_{\theta}(c(x,t)) &= \theta \varphi'(c(x,t)) - \theta_c c(x,t) \text{ to obtain} \\ \int_{\Omega} f'_{\theta}(c(t))^2 \, dx + \int_{\Omega} \underbrace{f''_{\theta}(c(t))}_{\geq -\theta_c} |\nabla c(t)|^2 \, dx \leq C \|\mu(t)\|_2 \end{split}$$

uniformly in $\theta > 0$.

by

A priori Estimates for c

 W_r^2 -estimate for c: Formally multiply

$$\mu(x,t) = -\Delta c(x,t) + f'_{\theta}(c(x,t))$$

by $f_{ heta}'(c(x,t)) = heta arphi'(c(x,t)) - heta_c c(x,t)$ to obtain

$$\int_{\Omega} f_{\theta}'(c(t))^2 dx + \int_{\Omega} \underbrace{f_{\theta}''(c(t))}_{\geq -\theta_c} |\nabla c(t)|^2 dx \leq C \|\mu(t)\|_2$$

uniformly in $\theta > 0$. Similarly, for $2 \le r < \infty$

$$\begin{split} \|f_{\theta}'(c(t))\|_{r} + \|c(t)\|_{W^{2}_{r}} &\leq C_{r} \left(\|\mu(t)\|_{r} + \|\nabla c(t)\|_{2}\right). \\ \\ \Rightarrow c \in L^{2}_{\mathsf{uloc}}([0,\infty); W^{2}_{6}(\Omega)) \end{split}$$

where

$$\|c\|_{L^2_{\mathrm{uloc}}([0,\infty);X)} = \sup_{t\geq 0} \|c\|_{L^2(t,t+1;X)}.$$

Modifications: Higher regularity in time in Besov spaces.

Higher Time Regularity for *c* $L^{\infty}(0,\infty; H_{(0)}^{-1})$ -estimate of $\partial_t c$: Multiplying $\partial_t^2 c + \Delta(\Delta c - \underbrace{f''(c)}_{\geq -\theta_c} \partial_t c) = -\partial_t (v \cdot \nabla c)$

by $-\Delta_N^{-1}\partial_t c$ yields

$$\begin{split} \|\partial_t c\|_{L^{\infty}(0,\infty;H_{(0)}^{-1})} + \|\nabla\partial_t c\|_{L^2(Q)} &\leq C(c_0) \left(1 + \|\partial_t v\|_{L^{\frac{4}{3}}_{uloc}(0,\infty;V'_n)}\right) \\ \text{where } V_n(\Omega) &= \{\varphi \in H^1(\Omega)^n : n \cdot \varphi|_{\partial\Omega} = 0\}. \\ \Rightarrow \mu \in L^{\infty}(0,\infty;H^1(\Omega)) \\ \Rightarrow c \in L^{\infty}(0,\infty;W^2_r(\Omega)), \ r = 6 \text{ if } d = 3 \text{ and } 1 < r < \infty \text{ if } d = 2. \end{split}$$

Higher Time Regularity for c $L^{\infty}(0,\infty; H_{(0)}^{-1})$ -estimate of $\partial_t c$: Multiplying $\partial_t^2 c + \Delta (\Delta c - \underbrace{f''(c)}_{\geq - heta_c} \partial_t c) = -\partial_t (v \cdot
abla c)$

by $-\Delta_N^{-1}\partial_t c$ yields

$$\begin{split} \|\partial_t c\|_{L^{\infty}(0,\infty;H_{(0)}^{-1})} + \|\nabla\partial_t c\|_{L^2(Q)} &\leq C(c_0) \left(1 + \|\partial_t v\|_{L^{\frac{4}{3}}_{uloc}(0,\infty;V'_n)}\right) \\ \text{where } V_n(\Omega) &= \{\varphi \in H^1(\Omega)^n : n \cdot \varphi|_{\partial\Omega} = 0\}. \\ \Rightarrow \mu \in L^{\infty}(0,\infty;H^1(\Omega)) \\ \Rightarrow c \in L^{\infty}(0,\infty;W_r^2(\Omega)), \ r = 6 \text{ if } d = 3 \text{ and } 1 < r < \infty \text{ if } d = 2. \\ \text{Problem: In general } \partial_t v \in L^{\frac{4}{3}}_{uloc}(0,\infty;H^{-1}(\Omega)^n) \not\subseteq L^{\frac{4}{3}}_{uloc}(0,\infty;V'_n)! \\ \text{Solution: Replace } \partial_t c \text{ by } h^{-\tau}\Delta_h c. \text{ Use } v \in B^{\tau}_{\frac{4}{3}\infty;uloc}([0,\infty);H^{-s}(\Omega)) \\ \text{with } 0 < s < \frac{1}{2}, \ \tau > \frac{2}{3} \text{ as well as } H^s_0(\Omega) = H^s(\Omega) \text{ and } H^{-s}(\Omega) = H^s(\Omega)' \\ \dots \Rightarrow c \in BUC([0,\infty);W^1_q(\Omega)), q > 3. \end{split}$$

Theorem (Double Obstacle Limit, A.'09)

There a subsequence of $(v^{\theta}, c^{\theta}, \mu^{\theta})_{\theta \in (0,1)}$ converges to (v^0, c^0, μ^0) solving

$$\partial_t v + v \cdot \nabla v - \operatorname{div}(\nu(c)Dv) + \nabla p = \mu_0 \nabla c \qquad \text{in } \Omega \times (0,\infty) \quad (7)$$

$$\operatorname{div} v = 0 \qquad \text{in } \Omega \times (0,\infty) \quad (8)$$

$$\partial_t c + v \cdot \nabla c = m \Delta \mu$$
 in $\Omega \times (0, \infty)$ (9)

$$b := \mu + \varepsilon \Delta c + \varepsilon^{-1} \theta_c c \in \partial I_{[-1,1]}(c) \quad in \ \Omega \times (0,\infty) \quad (10)$$

and $c(x,t) \in [-1,1]$ for all $(x,t) \in \Omega \times (0,\infty)$. Moreover,

$$egin{aligned} &(v^0,
abla c^0)\in L^\infty(0,\infty;L^2(\Omega)), &(
abla v^0,
abla \mu^0)\in L^2(0,\infty;L^2(\Omega)),\ &
abla^2c^0,b\in L^2_{loc}([0,\infty);L^6(\Omega)),\,c^0\in BUC([0,\infty);\mathcal{W}^1_q(\Omega)),q>c^0, \end{aligned}$$

For (v_0, c_0) sufficiently smooth:

- If d = 2, then the weak solution is unique and regular.
- ② If d = 3, there are some $0 < T_0 < T_1 < \infty$ such that the weak solution is regular and (locally) unique on $(0, T_0)$ and $[T_1, \infty)$.

Helmut Abels (U Regensburg)

Double Obstacle Limit

Open Questions

Question: How much does $\theta \ge 0$ influences the Ostwald ripening effect (for fixed ε)?

Open Questions

Question: How much does $\theta \ge 0$ influences the Ostwald ripening effect (for fixed ε)?

Simulation by S. Bartels

Question: Does (v(t), c(t)) converges to stationary solution as $t \to \infty$ if $\theta = 0$? (Known for $\theta > 0$, A. '07/'09)