

Stochastic Navier-Stokes-Coriolis Equations

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Contents

- ▶ Navier-Stokes-Coriolis equations
- ▶ Stochastic Navier-Stokes-Coriolis equations
(joint with M. Hieber (TU Darmstadt))
- ▶ Invariant measures
- ▶ Results in 2D
(partly joint with A. Es Sarhir (TU Berlin))

Navier-Stokes equation with Coriolis term

unbounded layer $\mathbb{R}^2 \times (0, b)$, $b > 0$

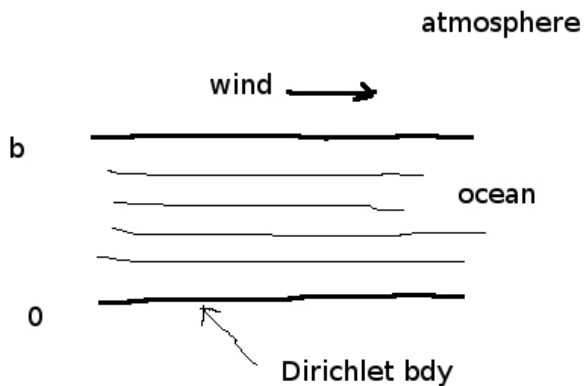
$$\left\{ \begin{array}{l} \partial_t u - \nu \Delta u + 2\omega(e_3 \times u) + (u \cdot \nabla)u + \nabla p = 0 \\ \operatorname{div} u = 0 \\ u(t, x_1, x_2, 0) = 0 \\ u(t, x_1, x_2, b) = u_b e_1 \\ u(0, \cdot) = u_0 \end{array} \right. \quad (1)$$

where

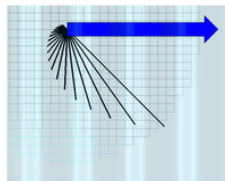
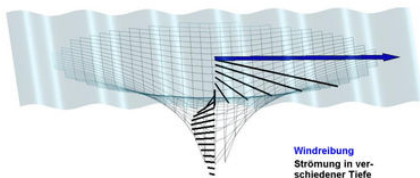
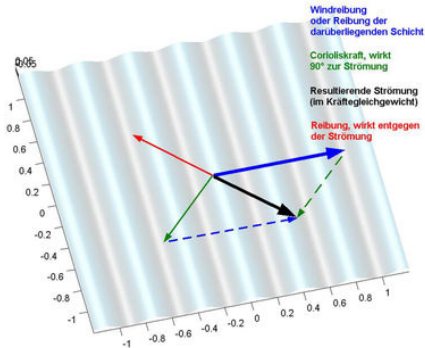
- ▶ u - velocity
- ▶ p - pressure
- ▶ $\nu > 0$ viscosity
- ▶ $\omega \in \mathbb{R}$ angular velocity
- ▶ $u_b e_1$, $u_b \in \mathbb{R}$ surface velocity

describes motion of rotating fluids under the influence of the Coriolis force

Background



Ekman spiral



(from Wikipedia)

Ekman spiral - analytic expression

explicit stationary solution of (1)

$$u_b^E(x_1, x_2, x_3) = \tilde{u}_b \begin{pmatrix} 1 - e^{-x_3/\delta} \cos(x_3/\delta) \\ e^{-x_3/\delta} \sin(x_3/\delta) \\ 0 \end{pmatrix}$$

$$p_b^E(x_1, x_2, x_3) = -\omega \tilde{u}_b x_2$$

with $\delta = \frac{b}{k\pi}$, $k \in \mathbb{Z}$, and

$$\tilde{u}_b = \begin{cases} u_b(1 - e^{-\frac{b}{\delta}})^{-1} & \text{if } k \text{ is even} \\ u_b(1 + e^{-\frac{b}{\delta}})^{-1} & \text{if } k \text{ is odd.} \end{cases}$$

Background: mathematical motivation

- ▶ (1) has a limiting equation for $|\omega| \rightarrow \infty$
- ▶ for arbitrary initial data u_0 and $|\omega| \gg 1$
 - ▶ ex. of global strong solutions
(Babin, Mahalov, Nicolaenko 1999)
 - ▶ ex. and uniqueness of mild solutions
(e.g.: Gallay, et al.: Global existence and Long-Time Asymptotics for Rotating Fluids in a 3D-Layer, 2009 on arXiv)

Known results: existence, stability

Function spaces

for $v = u - u_b^E$

$$\mathcal{D} := \left\{ v \in C_0^\infty(\mathbb{R}^2 \times (0, b))^3 : \operatorname{div} v = 0, v(x_1, x_2, 0) = 0 = v(x_1, x_2, b) \right\}$$

$$H := \overline{\mathcal{D}}^{L^2}(\mathbb{R}^2 \times (0, b))^3 \quad \text{and} \quad V := \overline{\mathcal{D}}^{H^{1,2}}(\mathbb{R}^2 \times (0, b))^3$$

Known results, ctd.

M. Hess (TU Darmstadt) 2009

- ▶ existence of global weak solutions

$$u = v + u_b^E, \quad v \in L^\infty(0, T; H) \cap L^2(0, T; V)$$

- ▶ if

$$\lambda_0 := \frac{\pi^2}{b^2} \left(\nu - \sqrt{2} \tilde{u}_b \left(\delta - (b + \delta) e^{-\frac{b}{\delta}} \right) \right) \geq 0$$

then \exists (at least one) global weak solution $u = v + u_b^E$ such that $\lim_{t \rightarrow \infty} \|v(t)\|_H = 0$, in addition

$$\|v(t)\|_H \leq e^{-\lambda t} \|v_0\|_H \text{ for any } \lambda < \lambda_0$$

Next step

- ▶ stability of u_b^E for the perturbed NSC-equation
- ▶ fluctuations around u_b^E for mean zero perturbations

Stochastic Navier-Stokes-Coriolis equations

restricting to (x_1, x_2) -periodic solutions u

$\Omega_b := \mathbb{T}^2 \times (0, b)$, \mathbb{T}^2 2D-Torus

Function spaces

for $v = u - u_b^E$

$$\mathcal{D} := \left\{ v \in C^\infty(\Omega_b)^3 : \operatorname{div} v = 0, v(x_1, x_2, 0) = 0 = v(x_1, x_2, b) \right\}$$

$$H := \overline{\mathcal{D}}^{L^2(\Omega_b)^3} \quad \text{and} \quad V := \overline{\mathcal{D}}^{H^{1,2}(\Omega_b)^3}$$

$$\left\{ \begin{array}{l} du_t = [\nu \Delta u_t - \omega(e_3 \times u_t) - (u_t \cdot \nabla) u_t + \nabla p_t] dt + dW_t \\ \operatorname{div} u_t = 0 \\ u_t(x_1, x_2, 0) = 0 \\ u_t(x_1, x_2, b) = u_b e_1 \end{array} \right. \quad (2)$$

$(W_t)_{t \geq 0}$ Wiener process on H

The Wiener process $(W_t)_{t \geq 0}$

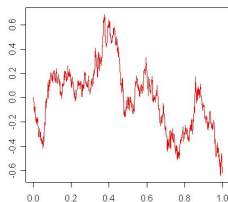
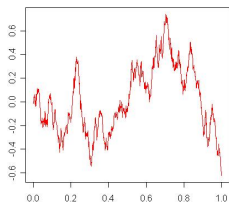
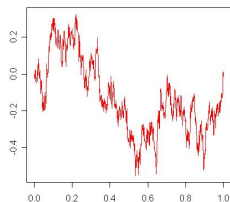
simplest example:

$$W_t = \sum_k \alpha_k \beta_k(t) e_k$$

where

- ▶ $\{e_k\}$ - complete orthonormal system of H
- ▶ $\{\alpha_k\} \subset \ell^2(\mathbb{R})$
- ▶ $\{\beta_k\}$ independent 1-dim Brownian motions on (Ω, \mathcal{F}, P)

Brownian motion β_k



cont. stoch. process, such that for $0 \leq t_0 < t_1 < \dots < t_n$ the increments

$$\beta_k(t_i) - \beta_k(t_{i-1}) \quad i = 1, \dots, n,$$

- ▶ independent
- ▶ $N(0, t_i - t_{i-1})$ -distributed

fundamental for the Ito-calculus: a typical path $t \mapsto \beta_k(t)(\omega)$ is

- ▶ of unbounded variation
- ▶ but of finite quadratic variation

$$\langle \beta_k(\omega) \rangle_t := \lim_{n \rightarrow \infty} \sum_{t_i \in \tau_n} (\beta_k(t_{i+1})(\omega) - \beta_k(t_i)(\omega))^2 = t$$

Stochastic evolution equation

orthogonal-projection $\Pi : L^2(\Omega_b)^3 \rightarrow H$, applied to (6), yields

$$du_t = [\nu A_S u_t - \omega \Pi(e_3 \times u_t) - \Pi(u_t \cdot \nabla) u_t] dt + dW_t \quad (3)$$

on the space H

$$A_S u = \nu \Pi \Delta u \quad \text{Stokes operator}$$

in the following: decompose

$$u_t = v_t + u_b^E \quad \text{w.r.t. the Ekman spiral } u_b^E$$

then v is a solution of

$$dv_t = [(\nu A_S + B)v_t - \Pi(v_t \cdot \nabla)v_t] dt + dW_t$$

where

$$B^{u_b^E} v := -\omega \Pi(e_3 \times v) - \Pi(u_b^E \cdot \nabla v) - \Pi(v_3 \partial_3 u_b^E), v \in D(A_S),$$

Perturbed stoch. Navier-Stokes eq

$$dv_t = [(\nu A_S + B)v_t - \Pi(v_t \cdot \nabla)v_t] dt + dW_t \quad (4)$$

Assumptions

A.1 $B : V \rightarrow H$, bounded, linear, s.th. $\exists \omega_0 > 0, \omega_1 \geq 0$ with

$$\langle (\nu A_S + B)u, u \rangle_H \leq -\omega_0 \|u\|_V^2 + \omega_1 \|u\|_H^2 \quad \forall u \in D(A_S).$$

A.2 $(W_t)_{t \geq 0}$ is a Wiener-process on H with covariance Q having finite trace

Rem A.1 implies that $(\nu A_S + B, D(A_S))$ generates an analytic C_0 -semigroup $(T_t)_{t \geq 0}$ satisfying $\|T_t\|_{op} \leq e^{(\omega_1 - \tilde{\omega}_0)t}, t \geq 0, \tilde{\omega}_0 = \omega_0 \frac{\pi^2}{b^2}$

Weak martingale solutions (up to time T)

given initial distribution μ_0 on H

a probability measure P on $\Omega = C([0, T]; H)$ such that the *canonical process* $v_t : \Omega \rightarrow H$, $\omega \mapsto \omega(t)$ satisfies



$$P \left[\sup_{t \in [0, T]} \|v_t\|_H + \int_0^T \|v_s\|_V^2 ds < \infty \right] = 1$$



$$\begin{aligned} M_t^\varphi &:= \langle v_t, \varphi \rangle_H \\ &\quad - \int_0^t \nu \langle v_s, A_S \varphi \rangle_H - \langle Bv_s, \varphi \rangle_H - \langle (v_s \nabla) \varphi, v_s \rangle_H ds \end{aligned}$$

is a continuous square-integrable martingale with

$$\langle M^\varphi \rangle_t = \|\sqrt{Q}\varphi\|_H^2 \cdot t \quad \text{for all } \varphi \in \mathcal{D}$$

▶ $P \circ v_0^{-1} = \mu_0$

Main existence result

Theorem (S., Hieber 2009) Let $\xi : \Omega \rightarrow H$ be \mathcal{F}_0 -measurable, in $\mathcal{L}^2(P)$, independent of $(W_t)_{t \geq 0}$.

Then \exists weak martingale solution $((v_t), (\tilde{W}_t))$ of (4) satisfying $\tilde{P} \circ v_0^{-1} = P \circ \xi^{-1}$

Moreover

$$\tilde{E} \left(\sup_{0 \leq t \leq T} \|v_t\|_H^2 + \int_0^T \|v_t\|_V^2 dt \right) < \infty$$

proof similar to the unperturbed case (e.g. Flandoli, Gatarek [PTRF 1995]) via Galerkin approximation, tightness, ...

Corollary $B = B_b^E$, then $u_t := v_t + u_b^E$ is a weak martingale solution of (3).

stationary solutions

a stochastic process $(v_t)_{t \geq 0}$ is called **stationary** if

$$\tilde{P} \circ v_{t+}^{-1} = \tilde{P} \circ v_t^{-1} \quad t \geq 0$$

in this case

$$\mu = \tilde{P} \circ v_t^{-1} \quad \text{is time-invariant}$$

μ describes the long-time fluctuations/statistics of $(v_t)_{t \geq 0}$

stationary martingale solutions of (4)

consider again the perturbed stochastic Navier-Stokes equation (4)

assume

$$\langle (\nu A_S + B)u, u \rangle_H \leq -\tilde{\omega}_0 \|u\|_V^2 \quad \forall u \in D(A_S)$$

with $\tilde{\omega}_0 := \omega_0 - \frac{1}{\nu} \frac{b^2}{\pi^2} \omega_1 > 0$

Theorem (S., Hieber, 2009) \exists a stat. martingale solution of (4)
Moreover, the invariant distribution $\mu = \tilde{P} \circ v_t^{-1}$ satisfies the moment estimate

$$\int (1 + \|v\|_V^2) e^{\varepsilon \|v\|_H^2} \mu(dv) < \infty \quad \text{for } \varepsilon < \varepsilon_0 := \frac{\tilde{\omega}_0}{\|Q\|_{op}}$$

open dependence on ω

at least: exponential moments ind. of ω

Navier-Stokes-Coriolis eq on \mathbb{T}^2

appropriate projection of (1) onto (x_1, x_2) -plane

$$\begin{cases} \partial_t u - \nu \Delta u + \ell e_3 \times u - (u \cdot \nabla) u + \nabla p = 0 \\ \operatorname{div} u = 0 \\ u(0, \cdot) = u_0 \end{cases} \quad (5)$$

where

▶ $\ell = \omega + \beta x_2$

▶

$$\ell e_3 \times u = \omega \begin{bmatrix} -u_2 \\ u_1 \end{bmatrix} + \beta x_2 \begin{bmatrix} -u_2 \\ u_1 \end{bmatrix}$$

ex. & uniq. well-known

Stoch. Navier-Stokes-Coriolis eq on \mathbb{T}^2

$$\begin{cases} du_t = [\nu \Delta u_t - \ell e_3 \times u_t - (u_t \cdot \nabla) u_t + \nabla p_t] dt + dW_t \\ \operatorname{div} u_t = 0 \end{cases} \quad (6)$$

$(W_t)_{t \geq 0}$ Wiener process on $H := \overline{\mathcal{D}}^{L^2(\mathbb{T}^2)^2}$

where

$$\mathcal{D} := \left\{ u \in C^\infty(\mathbb{T}^2)^2 : \operatorname{div} u = 0, \int_{\mathbb{T}^2} u \, dx = 0 \right\}$$

Stoch. Navier-Stokes-Coriolis eq on \mathbb{T}^2 , ctd.

applying the Helmholtz/Leray-projection yields

$$du_t = [\nu A_S u_t + Bu_t - \Pi(u_t \cdot \nabla) u_t] dt + dW_t \quad (7)$$

where

$$Bu = \Pi \left(\omega \begin{bmatrix} -u_2 \\ u_1 \end{bmatrix} + \beta x_2 \begin{bmatrix} -u_2 \\ u_1 \end{bmatrix} \right) = \beta \Pi \left(x_2 \begin{bmatrix} -u_2 \\ u_1 \end{bmatrix} \right)$$

Main features of B

- ▶ $\langle Bu, u \rangle = 0$ (cons. of energy)
- ▶ $\langle Bu, A_S u \rangle = 0$ (cons. of enstrophy)

Rem ex. of (stationary) weak martingale solution (should be) similar to 3D
(not done)

Direct construction of invariant measures

characterization via the ass. Kolmogorov operator

$$LF(u) := \frac{1}{2} \operatorname{tr}(QD^2F(u)) + \langle \nu A_S u + Bu - \Pi(u \cdot \nabla)u, DF(u) \rangle$$

$$F \in \mathcal{FC}_b^2 := \{F(u) = \varphi(\langle e_1, u \rangle, \dots, \langle e_n, u \rangle) : n \geq 1, \\ e_1, \dots, e_n \in D(A_S), \varphi \in C_b^2(\mathbb{R}^n)\}$$

Fact μ invariant for (7) implies μ infinitesimally invariant for L ,
i.e.,

$$L(\mathcal{FC}_b^2) \subset L^1(\mu) \quad \text{and} \quad \int LF d\mu = 0 \quad \forall F \in \mathcal{FC}_b^2$$

Invariant measures - finite trace

Theorem (S. 2009)

- ▶ $\text{tr}_H(Q) < \infty \implies \exists \mu$ inf. invariant for L with

$$\int (1 + \|u\|_V^2) e^{\varepsilon \|u\|_H^2} \mu(du) < \infty$$

$$\text{if } \varepsilon < \frac{\nu}{\|Q\|_{L(H)}}$$

- ▶ $\text{tr}_V(Q) < \infty \implies \mu$ satisfies in addition

$$\int (1 + \|u\|_{H^2}^2) e^{\varepsilon \|u\|_V^2} \mu(du) < \infty \quad (8)$$

$$\text{if } \varepsilon < \frac{\nu}{\|Q\|_{L(V)}}$$

L^p -uniqueness - finite trace

Theorem (S. 2009)

Let μ be inf. invariant, satisfying (8), then (L, \mathcal{FC}_b^2) is L^p -unique for all $p \in [1, \infty[$

basic ingredient: pointwise gradient estimate for finite-dim. Galerkin approx.
 L^N of L

$$\|DT_t^N F(u)\|_H \leq \exp\left(\frac{\varepsilon\nu}{4(\nu^2 - \varepsilon\|\sqrt{Q}\|_{L(H^2, \nu)}^2)}\|u\|_V^2 + C(\varepsilon, Q, \nu)t\right)\|F\|_{Lip(H)}$$

uniformly in N

Rem $B = 0$ result obtained by Barbu, Da Prato (2004)

white noise $Q = id$

recall invariants for $\Pi(u \cdot \nabla)u$ and B

- ▶ energy $\langle \Pi(u \cdot \nabla)u, u \rangle = 0 = \langle Bu, u \rangle$
- ▶ enstrophy $\langle \Pi(u \cdot \nabla)u, A_S u \rangle = 0 = \langle Bu, A_S u \rangle$

implies

$$L_{\sigma, \nu} F(z) = \frac{\sigma^2}{2} \text{tr} (D^2 F)(u) + \langle \nu A_S u + Bu - \Pi(u \cdot \nabla)u, DF(u) \rangle$$

has (inf.) invariant measure $\mu_{2, \sigma, \nu} = N(0, \frac{\sigma^2}{\nu} (-A_S)^{-1})$

(independent of ω)

A priori estimate in the case $C = 0$

Theorem (S. 2007, IDAQP) Let $\alpha \in (-1, 0)$. Then there exists $C(\alpha)$ such that:

- ▶ If $\nu^3 > \sigma^2 C(\alpha)$ for some $\alpha \in (-1, 0)$ then the closure $\bar{L}_{\sigma, \nu}$ generates a Markovian C_0 -semigroup on $L^1(\mu_{2, \sigma, \nu})$.
- ▶ The associated resolvent $\bar{R}_{\sigma, \nu, \lambda} = (\lambda - \bar{L}_{\sigma, \nu})^{-1}$, $\lambda > 0$, satisfies

$$\int |D\bar{R}_{\sigma, \nu, \lambda} F(u)|_{H^{2+\alpha}}^2 \mu(du) \leq \frac{\nu^2}{\lambda(\nu^3 - \sigma^2 C(\alpha))} \int |DF(u)|_{H^{1+\alpha}}^2 \mu(du)$$

in part. $\bar{R}_{\sigma, \nu, \lambda}$ operates as a bounded operator

$$\bar{R}_{\sigma, \nu, \lambda} : H_{H^{1+\alpha}}^{1,2}(\mu_{2, \sigma, \nu}) \rightarrow H_{H^{2+\alpha}}^{1,2}(\mu_{2, \sigma, \nu})$$

for $\alpha \in (-1, 0)$.

Consequences

- ▶ existence of diffusions generated by $\bar{L}_{\sigma, \nu}$ using Dirichlet form techniques
- ▶ uniqueness of ass. martingale problem

A priori estimate in the case $C = 0$

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$$\int |D\bar{R}_{\sigma, \nu, \lambda} F(u)|_{H^{2+\alpha}}^2 \mu(du) \leq \frac{\nu^2}{\lambda(\nu^3 - \sigma^2 C(\alpha))} \int |DF(u)|_{H^{1+\alpha}}^2 \mu(du).$$

$B \neq 0$ **conjecture** for ν^3/σ^2 sufficiently large

- ▶ the same integr. gradient estimates hold
- ▶ $(L_{\sigma, \nu}, \mathcal{F}C_b^2)$ is L^1 -unique

NB

- ▶ invariant measure is independent of angular velocity ω
- ▶ can therefore study the limit $|\omega| \rightarrow \infty$ in $L^1(\mu_{2, \sigma, \nu})$

rough noise $Q = (-A_S)^{-\frac{1}{2}-\varepsilon}$

ex. of (stationary) weak martingale solution (should be) similar to the case without Coriolis forcing term (not done)

ex. proof for the stoch. Navier-Stokes without Coriolis forcing term by Flandoli, Gatarek [PTRF, 1995] suggests

$$\int \log(1 + \|u\|^2) \mu(du) < \infty$$

for an invariant distribution

can be improved to

$$\implies \int \left(1 + \|u\|_\gamma^2\right) \|u\|_H^{2m} \mu(du) < \infty \quad \forall m \geq 0, \gamma < \frac{1}{4} + \varepsilon$$

cf. **Es-Sarhir, S., 2009**, (to appear in JFA) Improved moment estimates for inv. measures of semilinear diffusions in Hilbert spaces and applications

improved moment estimates

improved moment estimates for spde of the type

$$dX_t = (AX_t + B(X_t)) dt + \sqrt{Q} dW_t \text{ on } H$$

A self-adjoint, of negative type, compact resolvent
satisfying

$$\begin{aligned} \langle Ay + B(y + w), y \rangle &\leq -\alpha_1 \|y\|_{V_{\frac{1}{2}}}^2 + \alpha_2 \|w\|_{V_{\gamma_0}}^s \\ &\quad + \alpha_3 \|w\|_{V_{\gamma_0}}^s \|y\|_{V_{\gamma_1}}^2 + \alpha_4 \end{aligned}$$

Rem $\alpha_3 = 0$: exponential moments (cf. Es Sarhir, S., JEE 2008)

$\alpha_3 > 0$ - main ingredient

Prop (Da Prato, Debussche, ...) Let $\delta \in]0, \frac{1}{2}[$ and $\gamma \in \mathbb{R}$. Then

$$\sup_{0 \leq t \leq T} \|W_{A-\lambda}(t)\|_{V_\gamma}^2 \leq C_\delta^2 \sum_{k=1}^{\infty} \frac{\lambda_k^{2\gamma}}{(\lambda + \lambda_k)^{2\delta}} M_k(\delta, T)^2$$

where

▶ $W_{A-\lambda}(t) = \int_0^t e^{(t-s)A} \sqrt{Q} dW_s$

▶ (λ_k) - eigenvalues of $-A$

▶

$$M_k(\delta, T) := \sup_{0 \leq s < t \leq T} \frac{|\beta_k(t) - \beta_k(s)|}{|t - s|^\delta}$$

ind. r.v., having finite moments of any order. Here,

$$\beta_k(t) := \langle W(t), e_k \rangle \quad (e_k)_k \text{ eigenfunctions of } A$$

ind. 1-dim Brownian motions

in part.: if for some $\varepsilon > 0$

$$\sum_{k=1}^{\infty} \lambda_k^{-2(\delta-\gamma-\varepsilon)} < \infty$$

then

$$\sup_{0 \leq t \leq T} \|W_{A-\lambda}(t)\|_{V_\gamma}^2 \leq \lambda^{-2\varepsilon} M_{\delta,\gamma,\varepsilon}$$

where $M_{\delta,\gamma,\varepsilon}$ ind. of λ , finite moments