Qualitative Behavior of Solutions for the Two-Phase Navier-Stokes Equations with Surface Tension

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Notations

 $\Omega \subset \mathbb{R}^n$ open bounded domain with smooth boundary, $\Omega_i(t)$ subdomains occupied by immiscible incompressible fluid i = 1, 2, $\Gamma(t)$ interface separating the two phases. No boundary contact, i.e $\Gamma(t) \cap \partial \Omega = \emptyset$, no phase transitions, isothermal conditions and no external forces.

$$\begin{split} & u = u(t,x) \text{ velocity field, } \pi = \pi(t,x) \text{ pressure field} \\ & S(t,x) \text{ stress tensor} \\ & E(t,x) := \frac{1}{2} (\nabla u(t,x) + \nabla u(t,x)^{\mathsf{T}}) \text{ rate of strain tensor} \\ & \rho_i > 0 \text{ densities, } \mu_i > 0 \text{ viscosities in the phases} \\ & \nu_{\mathsf{\Gamma}}(t,x) \text{ the normal at } x \in \mathsf{\Gamma}(t) \text{ directed into } \Omega_2(t) \\ & V_{\mathsf{\Gamma}}(t,x) = (u(t,x)|\nu_{\mathsf{\Gamma}}(t,x)) \text{ normal velocity of } \mathsf{\Gamma}(t) \\ & H_{\mathsf{\Gamma}}(t,x) = -\text{div}_{\mathsf{\Gamma}}\nu_{\mathsf{\Gamma}}(t,x) \text{ curvature of } \mathsf{\Gamma}(t) \end{split}$$

$$\llbracket \phi \rrbracket = \lim_{h \to 0+} [\phi(t, x + h\nu_{\Gamma}(t, x)) - \phi(t, x - h\nu_{\Gamma}(t, x))]$$

is the jump of the quantity ϕ accross $\Gamma(t)$.

The Two-Phase Navier-Stokes Problem

In the bulk phases:

$$\begin{array}{rcl} \partial_t(\rho u) + \nabla \cdot (\rho u \otimes u - S) &=& 0, & x \in \Omega \backslash \Gamma(t), \ t > 0, \\ \nabla \cdot u &=& 0, & x \in \Omega \backslash \Gamma(t), \ t > 0, \\ \mu(\nabla u + \nabla u^{\mathsf{T}}) - \pi I &=& S, & x \in \Omega \backslash \Gamma(t), \ t > 0. \end{array}$$

At the interface:

Initial conditions:

$$u(0,x) = u_0(x)$$
 $x \in \Omega$, $\Gamma(0) = \Gamma_0$.

No-slip boundary conditions at $\partial \Omega$, i.e. u = 0 on $\partial \Omega$.

One-phase problems with surface tension: BEALE 1980, 1984, SOLONNIKOV 1987-, TANI & TANNAKA 1996, SHIMIZU & SHIBATA 2005-.

Two phase problems with surface tension: Denisova 1988-, Tani 1996, Tannaka 1995, Shimizu & Shibata 2009, Prüss & Simonett 2009.

Of particular relevance for convergence of solutions are the papers of BEALE and SOLONNIKOV, both in the *one-phase case*.

Here we are interested in the **qualitative and asymptotic behaviour** of solutions of the two-phase problem.

Fix p > n + 2, let $\partial \Omega \in C^3$, and suppose

$$\Gamma_0 \in W_p^{3-2/p}, \quad u_0 \in W_p^{2-2/p}(\Omega \setminus \Gamma_0)^n.$$

Assume the compatibility conditions

div $u_0 = 0$ in $\Omega \setminus \Gamma_0$, u = 0 on $\partial \Omega$, $\llbracket \mathcal{P}_{\Gamma_0} \mu \mathcal{E}_0 \nu_{\Gamma_0} \rrbracket = 0$, $\llbracket u_0 \rrbracket = 0$ on Γ_0 ,

where $E_0 = \frac{1}{2}(\nabla u_0 + \nabla u_0^T)$, and $\mathcal{P}_{\Gamma_0} = I - \nu_{\Gamma_0} \times \nu_{\Gamma_0}$.

Then there exists $t_0 = t_0(u_0, \Gamma_0) > 0$ and a unique classical solution (u, π, Γ) of the problem on $(0, t_0)$. The set

 $\bigcup_{t\in(0,t_0)}\{t\}\times \Gamma(t)$

is a real analytic manifold, and with

 $\mho := \{(t, x) \in (0, t_0) \times \Omega, \ x \notin \Gamma(t)\},\$

the function (u, π) : $\mho \to \mathbb{R}^{n+1}$ is real analytic.

Strategy for the proof

Approximate Γ_0 by a smooth hypersurface Σ .

a) Transformation to a Fixed Domain (Hanzawa) Let d(x) denote the signed distance of $x \in \mathbb{R}^n$ to Σ , and $\Pi(x)$ the projection of $x \in \mathbb{R}^n$ to Σ . Then

 $egin{aligned} & \Lambda: \Sigma imes (-a,a) o \mathbb{R}^n \ & \Lambda(p,r) := p + r
u_{\Sigma}(p), \quad \Lambda^{-1}(x) = (\Pi(x), d(x)) \end{aligned}$

is a diffeomorphism from $\Sigma \times (-a, a)$ onto $\mathcal{R}(\Lambda) = \{x \in \mathbb{R}^n : |d(x)| < a\}$, provided

 $0 < a < \min\{r, 1/\kappa_j(x) : j = 1, \dots, n-1, x \in \Sigma\},$

where $\kappa_j(x)$ mean the principal curvatures of Σ at $x \in \Sigma$ and

$$\overline{B}_r(p\pm r
u_{\Sigma}(p))\cap \Sigma=\{p\}, \quad p\in \Sigma.$$

Use this to parameterize $\Gamma(t)$ over Σ :

 $\Gamma(t): p\mapsto p+h(t,p)\nu_{\Sigma}(p), \quad p\in \Sigma, \ t\geq 0.$

Extend this diffeomorphism to all of Ω :

 $\Theta(t,x) = x + \chi(d(x))h(t,\Pi(x))\nu_{\Sigma}(\Pi(x)).$

Here χ denotes a suitable cut-off function. This way $\Omega \setminus \Gamma(t)$ is transformed to the fixed domain $\Omega \setminus \Sigma$. Set

 $ar{u}(t,x) = u(t, \Theta(t,x)), \ ar{\pi}(t,x) = \pi(t, \Theta(t,x)), \quad t>0, \ x\in \Omega ackslash \Sigma.$

This yields the problem (drop the bars!)

$$\begin{split} \rho \partial_t u &- \mu \mathcal{A}(h) u + \mathcal{G}(h) \pi = \mathcal{R}(u, h) & \text{ in } \Omega \setminus \Sigma, \\ (\mathcal{G}(h)|u) &= 0 & \text{ in } \Omega \setminus \Sigma, \\ u &= 0 & \text{ on } \partial \Omega, \\ \llbracket -\mu(\mathcal{G}(h)u + [\mathcal{G}(h)u]^{\mathsf{T}}) + \pi \rrbracket \nu_{\mathsf{\Gamma}}(h) &= \sigma H_{\mathsf{\Gamma}}(h) \nu_{\mathsf{\Gamma}}(h) \text{ on } \Sigma, \\ \llbracket u \rrbracket &= 0 & \text{ on } \Sigma, \\ \partial_t h - (u|\nu_{\Sigma}) &= -(u|\alpha(h)), & \text{ on } \Sigma, \\ u(0) &= u_0, \quad h(0) &= h_0. \end{split}$$

Here $\mathcal{A}(h)$ and $\mathcal{G}(h)$ denote the transformed Laplacian, resp. gradient. With the curvature tensor L_{Σ} and the surface gradient ∇_{Σ} we have

 $u_{\Gamma}(h) = \beta(h)(\nu_{\Sigma} - \alpha(h)), \qquad \alpha(h) = M(h)\nabla_{\Sigma}h,$ $M(h) = (I - hL_{\Sigma})^{-1}, \qquad \beta(h) = (1 + |\alpha(h)|^2)^{-1/2},$

and

$$V = (\partial_t \Theta | \nu_{\Gamma}) = \partial_t h(\nu_{\Gamma} | \nu_{\Sigma}) = \beta(h) \partial_t h.$$

The curvature $H_{\Gamma}(h)$ becomes

 $\begin{aligned} H_{\Gamma}(h) &= \beta(h) \{ \operatorname{tr}[M(h)(L_{\Sigma} + \nabla_{\Sigma} \alpha(h))] \\ &- \beta^{2}(h)(M(h)\alpha(h)|[\nabla_{\Sigma} \alpha(h)]\alpha(h)) \}, \end{aligned}$

a differential expression involving second order derivatives of h only linearly. Linearization of $H_{\Gamma}(h)$ at h = 0: $H'_{\Gamma}(0) = \operatorname{tr} L^2_{\Sigma} + \Delta_{\Sigma}$, where Δ_{Σ} denotes the Laplace-Beltrami operator on Σ .

Rewrite this problem in quasilinear form.

$$\rho \partial_t u - \mu \Delta u + \nabla \pi = F(u, \pi, h) \quad \text{in } \Omega \setminus \Sigma,$$

$$\operatorname{div} u = F_d(u, h) \quad \text{in } \Omega \setminus \Sigma,$$

$$u = 0 \quad \text{on } \partial \Omega,$$

$$\llbracket -\mu (\nabla u + \nabla u^{\mathsf{T}}) + \pi \rrbracket \nu_{\Sigma} - \sigma (\Delta_{\Sigma} h) \nu_{\Sigma} = G(u, \llbracket \pi \rrbracket, h) \quad \text{on } \Sigma,$$

$$\llbracket u \rrbracket = 0 \quad \text{on } \Sigma,$$

$$\partial_t h - (u | \nu_{\Sigma}) = G_h(u, h) \quad \text{on } \Sigma,$$

$$u(0) = u_0, \quad h(0) = h_0$$
(2)

The right hand sides in this problem consist of lower order terms and of terms of the same order appearing on the left, but carrying a factor $|\nabla_{\Sigma} h|$, which is small by construction.

b) The Linear Problem

Establish maximal L_p -regularity for the linear problem defined by the l.h.s. of (2). In particular, find the right spaces for the data, such that the *solution-to-data map* for the following inhomogeneous linear problem becomes an *isomorphism*.

> $\rho \partial_t v - \mu \Delta v + \nabla q = f_v \quad \text{in } \Omega \setminus \Sigma,$ $\operatorname{div} v = f_d \quad \operatorname{in } \Omega \setminus \Sigma$ $v = 0 \quad \text{on } \partial \Omega,$ $\llbracket -\mu (\nabla v + \nabla v^{\mathsf{T}}) + q \rrbracket v_{\Sigma} - \sigma (\Delta_{\Sigma} h) v_{\Sigma} = g \quad \text{on } \Sigma,$ $\llbracket v \rrbracket = 0 \quad \text{on } \Sigma,$ $\partial_t h - (v | v_{\Sigma}) = g_h \quad \text{on } \Sigma,$ $v(0) = v_0, \quad h(0) = h_0.$ (3)

This is proved by localization and perturbation, and a corresponding result for a flat interface, i.e. $\Omega = \mathbb{R}^n$, $\Sigma = \mathbb{R}^{n-1} \times \{0\}$; see PRÜSS & SIMONETT 2009.

The solutions of the transformed problem will belong to the following class:

 $v \in H^1_p(J; L_p(\Omega)^n) \cap L_p(J; H^2_p(\Omega \setminus \Sigma)^n), \quad q \in L_p(J; \dot{H}^1_p(\Omega \setminus \Sigma)),$

$$\llbracket q \rrbracket \in W^{1/2-1/2p}_p(J; L_p(\Sigma)) \cap L_p(J; W^{1-1/p}_p(\Sigma)),$$

 $h \in W_p^{2-1/2p}(J; L_p(\Sigma)) \cap H_p^1(J; W_p^{2-1/p}(\Sigma)) \cap L_p(J; W_p^{3-1/p}(\Sigma)),$ provided

$$v_0 \in W^{2-2/p}_p(\Omega \setminus \Sigma)^n, \quad h_0 \in W^{3-2/p}_p(\Sigma),$$

and the natural compatibility conditions hold.

Let $1 , <math>\rho_j$, μ_j , σ be positive constants, j = 1, 2; set J = [0, a]. Then the two-phase Stokes problem (3) admits a unique solution (v, q, h) with regularity

 $v \in H^1(J; L_p(\Omega))^n \cap L_p(J; H^2_p(\Omega \setminus \Sigma))^n, \quad q \in L_p(J; \dot{H}^1_p(\Omega \setminus \Sigma)),$

$$[\![q]\!] \in W^{1/2-1/2p}_{
ho}(J; L_{
ho}(\Sigma)) \cap L_{
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ho}(\Sigma)),$$

 $h \in W_p^{2-1/2p}(J; L_p(\Sigma)) \cap H_p^1(J; W_p^{2-1/p}(\Sigma) \cap L_p(J; W_p^{3-1/p}(\Sigma))).$ if and only if the data $(v_0, h_0, f, f_d, g, g_h)$ satisfy the following regularity and compatibility conditions.

(a)
$$f_{v} \in L_{p}(J \times \Omega)^{n}, v_{0} \in W_{p}^{2-2/p}(\Omega \setminus \Sigma)^{n}, v_{0} = 0 \text{ on } \partial\Omega;$$

(b) $f_{d} \in H_{p}^{1}(J; H_{p}^{-1}(\Omega \setminus \Sigma)) \cap L_{p}(J; H_{p}^{1}(\Omega \setminus \Sigma)), \quad \text{div } v_{0} = f_{d}(0);$
(c) $g \in W_{p}^{1/2-1/2p}(J; L_{p}(\Sigma)) \cap L_{p}(J; W_{p}^{1-1/p}(\Sigma));$
(d) $\llbracket v_{0} \rrbracket = 0, \quad \mathcal{P}_{\Sigma} \llbracket \mu(\nabla v_{0} + \nabla v_{0}^{T}) \rrbracket = g_{v}(0);$
(e) $h_{0} \in W_{p}^{3-2/p}(\Sigma);$
(f) $g_{h} \in W_{p}^{1-1/2p}(J; L_{p}(\Sigma)) \cap L_{p}(J; W_{p}^{2-1/p}(\Sigma)).$

The solution map $(v_0, h_0, f, f_d, g, g_h) \mapsto (v, q, h)$ is continuous between the corresponding spaces.

c) The Nonlinear Problem

Based on *maximal regularity*, use the *implicit function theorem* to obtain local well-posedness of the nonlinear problem.

Use a variant of ANGENENT'S *parameter trick* to obtain real analyticity via *maximal regularity* and the *implicit function theorem*; cf. ESCHER, PRÜSS & SIMONETT 2003.

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(c) $g \in W_{p}^{1/2-1/2p}(J; L_{p}(\Sigma)) \cap L_{p}(J; W_{p}^{1-1/p}(\Sigma));$
(d) $\llbracket v_{0} \rrbracket = 0, \quad \mathcal{P}_{\Sigma} \llbracket \mu(\nabla v_{0} + \nabla v_{0}^{T}) \rrbracket = g_{v}(0);$
(e) $h_{0} \in W_{p}^{3-2/p}(\Sigma);$
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The Induced Semiflow

Recall that the closed C^2 -hypersurfaces contained in Ω form a C^2 -manifold, which we denote by \mathcal{MH}^2 . Charts are obtained via parametrization over a fixed hypersurface. As an ambient space for the phase-manifold \mathcal{PM} of the two-phase Navier-Stokes problem with surface tension we consider the product space $C(\overline{\Omega})^n \times \mathcal{MH}^2$. We define \mathcal{PM} as follows.

 $\mathcal{PM} :=$

 $\{(u,\Gamma) \in C(\bar{\Omega})^n \times \mathcal{MH}^2 : u \in W_p^{2-2/p}(\Omega \setminus \Gamma)^n, \, \Gamma \in W_p^{3-2/p}, \\ \operatorname{div} u = 0 \text{ in } \Omega \setminus \Gamma, \, u = 0 \text{ on } \partial\Omega, \, \mathcal{P}_{\Gamma}\llbracket \mu E \rrbracket \nu_{\Gamma} = 0 \text{ on } \Gamma \}.$ (4)

The charts for this manifold are obtained by the charts induced by \mathcal{MH}^2 , followed by a HANZAWA transformation.

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The charts for this manifold are obtained by the charts induced by \mathcal{MH}^2 , followed by a HANZAWA transformation.

Observe that the compatibility conditions

div u = 0 in $\Omega \setminus \Gamma$, u = 0 on $\partial \Omega$, $\mathcal{P}_{\Gamma} \mu \llbracket (\nabla u + \nabla u^{\mathsf{T}}) \rrbracket \nu_{\Gamma} = 0$, $\llbracket u \rrbracket = 0$ on Γ ,

as well as regularity are preserved by the solutions. This yields the following result

Theorem

Let p > n + 2. Then the two-phase Navier-Stokes problem with surface tension generates a local semiflow on the phase-manifold \mathcal{PM} . Each solution (u, Γ) exists on a maximal time interval $[0, t_*)$.

The energy functional and equilibria

Define the energy functional by means of

$$\Phi(u,\Gamma) := \frac{1}{2} \|\rho^{1/2}u\|_{L_2(\Omega)}^2 + \sigma \operatorname{mes} \Gamma(t).$$

Then

$$\partial_t \Phi(u,\Gamma) + 2 \|\mu^{1/2} E\|_{L_2(\Omega)}^2 = 0,$$

hence the energy functional is a Ljapunov functional, even a strict one.

Proposition

Let $\rho_i, \mu_i, \sigma > 0$ be constants. Then (a) The **energy equality** is valid for smooth solutions.

(b) The equilibria are zero velocities, constant pressures in the

phase-components, the dispersed phase is a union of nonintersecting balls.

(c) The energy functional is a strict Ljapunov-functional.
(d) The critical points of the energy functional for constant phase volumes are precisely the equilibria.

The Stability Result

Assuming, for simplicity, that the phases are connected, we denote by

 $\mathcal{E} := \{(0, S_R(x_0)): x_0 \in \Omega, R > 0, S_R(x_0) \subset \Omega\}$

the set of equilibria without boundary contact. Note that \mathcal{E} forms a real analytic manifold of dimension n + 1.

Fix any such equilibrium $(0, \Sigma) \in \mathcal{E}$. We consider the behaviour of the solutions near this steady state.

Here we have to use the *full linearization* of the problem at an equilibrium $(0, \Sigma)$ i.e. at (u, h) = (0, 0), and for this reason we have to replace Δ_{Σ} in the linear problem (3) by

$$\mathcal{A}_{\Sigma} = \mathcal{H}_{\Gamma}'(0) = rac{n-1}{R^2} + \Delta_{\Sigma}.$$

It is well-known that \mathcal{A}_{Σ} is *selfadjoint*, *negative semidefinite* and has *compact resolvent* in $L_2(\Sigma)$; $\lambda_0 = 0$ is an eigenvalue of dimension *n*.

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The equilibrium $(0, \Sigma)$ is stable in the sense that for each $\varepsilon \in (0, \varepsilon_0]$ there exists $\delta(\varepsilon) > 0$ such that for all initial values (u_0, Γ_0) subject to

$$\operatorname{dist}_{W_p^{3-2/p}}(\Gamma_0, \Sigma) \leq \delta(\varepsilon) \quad \text{ and } \quad \|u_0\|_{W_p^{2-2/p}(\Omega \setminus \Gamma_0)} \leq \delta(\varepsilon)$$

there exists a unique global solution $(u(t), \Gamma(t))$ of the problem, and it satisfies

 $\operatorname{dist}_{W^{3-2/p}_p}(\Gamma(t),\Sigma) \leq \varepsilon \quad \text{ and } \quad \|u(t)\|_{W^{2-2/p}_p(\Omega\setminus\Gamma(t))} \leq \varepsilon, \ t\geq 0.$

Moreover, as $t \to \infty$ each of these solutions $(u(t), \Gamma(t))$ converges to a probably different equilibrium $(0, \Sigma_{\infty})$ in the same topology, i.e.

$$\lim_{t\to\infty} \left({\rm dist}_{W^{3-2/p}_{\rho}}(\Gamma(t), \Sigma_{\infty}) + \|u(t)\|_{W^{2-2/p}_{\rho}(\Omega\setminus\Gamma(t))} \right) = 0.$$

The convergence is at exponential rate.

Sketch of the Proof

As a base space we use

$$X_0 = L_{p,\sigma}(\Omega)^n \times W_p^{2-1/p}(\Sigma),$$

and we set

$$ar{X}_1 = W_p^2(\Omega \setminus \Sigma)^n imes W_p^{3-1/p}(\Sigma)$$

Define a closed linear operator in X_0 by means of

$$A(\mathbf{v}, \mathbf{h}) = (-(\mu/\rho)\Delta \mathbf{v} + \rho^{-1}\nabla q, -(\mathbf{v}|\nu_{\Sigma})),$$

with domain $X_1 := D(A) \subset \bar{X}_1$

$$\begin{split} D(A) &= \{ (v,h) \in \bar{X}_1 \cap X_0 : v = 0 \text{ on } \partial \Omega, \ \llbracket v \rrbracket = 0 \text{ and} \\ & \llbracket \mathcal{P}_{\Sigma} \mu (\nabla v + \nabla v^{\mathsf{T}}) \nu_{\Sigma} \rrbracket = 0 \text{ on } \Sigma \}. \end{split}$$

Here $q \in \dot{W}^1_p(\Omega \setminus \Sigma)$ is determined as the solution of the *weak* transmission problem

$$\begin{aligned} (\rho^{-1} \nabla q | \nabla \phi)_{L_2} &= ((\mu/\rho) \Delta v | \nabla \phi)_{L_2}, \quad \phi \in W^1_{\rho'}(\Omega), \\ \llbracket q \rrbracket &= \llbracket \mu ((\nabla v + \nabla v^{\mathsf{T}}) \nu_{\Sigma} | \nu_{\Sigma}) \rrbracket + \sigma \mathcal{A}_{\Sigma} h \quad \text{on } \Sigma. \end{aligned}$$

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Then with z = (v, h) and $f = (f_v, g_h)$ as well as $z_0 = (v_0, h_0)$, system (3) can be rewritten as the abstract evolution equation

$$\dot{z} + Az = f, \quad t > 0, \quad z(0) = z_0,$$
 (5)

provided $f_d = 0$ and g = 0.

Since (3) has maximal L_p -regularity, the abstract problem (5) has maximal L_p -regularity, as well. In particular, -A generates an analytic C_0 -semigroup in X_0 .

We can show the following properties:

(i) The set of equilibria *E* is an (n + 1)-dimensional smooth manifold.
(ii) The kernel N(A) is isomorphic to the tangent space T_(0,Σ)*E*.
(iii) N(A) ⊕ R(A) = X₀, i.e. the eigenvalue λ₀ = 0 of A is semi-simple.
(iv) σ(A) \ {0} ⊂ C₊, i.e. the semigroup e^{-At}_{|P(A)} is exponentially stable.

This allows for the use of the *generalized principle of linearized stability*, see PRÜSS, SIMONETT & ZACHER 2009, to prove the theorem.

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There are basically two obstructions against global existence:

- **regularity**: the norms of either u(t) or $\Gamma(t)$ become unbounded;
- **geometry**: the topology of the interface changes or the interface touches the boundary of Ω .

Note that the *phase volumes* are conserved by the semiflow!

We say that a solution (u, Γ) satisfies a uniform ball condition, if there is a radius r > 0 such that for each $t \in [0, t_*)$ and at every point $p \in \Gamma(t)$ we have

 $\overline{B}_r(p \pm r\nu_{\Gamma(t)}(p)) \cap \Gamma(t) = \{p\}.$

Combining the above results, we obtain the following theorem on the asymptotic behavior of solutions.

Theorem

Let p > n + 2. Suppose that (u, Γ) is a solution of the two-phase Navier-Stokes problem with surface tension on the maximal time interval $[0, t_*)$. Assume the following on $[0, t_*)$: (i) $||u(t)||_{W_p^{2-2/p}} + ||\Gamma(t)||_{W_p^{3-2/p}} \le M < \infty$; (ii) (u, Γ) satisfies a uniform ball condition. Then $t_* = \infty$, i.e. the solution exists globally, and it converges in \mathcal{PM} at exponential rate to an equilibrium $(0, \Gamma_\infty) \in \mathcal{E}$. The converse is also true.

The idea of the proof is as follows. Assuming (i) and (ii) we show that the solution is global. The energy is a strict Ljapunov functional, hence the limit set $\omega(u, \Gamma)$ of a solution is contained in the set \mathcal{E} of equilibria. A compactness argument shows $\omega(u, \Gamma) \neq \emptyset$, hence the solution comes close to \mathcal{E} . Then we may apply the local convergence result. Combining the above results, we obtain the following theorem on the asymptotic behavior of solutions.

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