Blow-up solutions and its singularities of the heat equations with a nonlinear boundary condition

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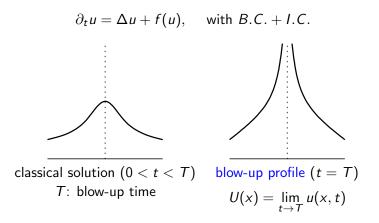
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Introduction

Typical blow-up phenomenon for semilinear parabolic equations



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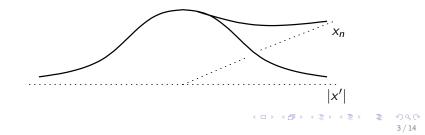
Our equation and assumption

We study blow-up phenomenon of positive solutions to

$$\begin{cases} \partial_t u = \Delta u & \text{in } \mathbb{R}^n_+ \times (0, T), \\ \partial_\nu u = u^q & \text{on } \partial \mathbb{R}^n_+ \times (0, T) & + \text{I.C.} \end{cases}$$
(P)

where $q \in (1, n/(n-2))$.

Definition. We call $u(x) \times_n$ -axial symmetric, if $u(x) = u(|x'|, x_n)$. Assumption (monotonicity). $x' \cdot \nabla' u_0(x) \le 0$, $\partial_n u_0(x) \le 0$.



Properties of equation (P)

$$\begin{cases} \partial_t u = \Delta u & \text{in } \mathbb{R}^n_+ \times (0, T), \\ \partial_\nu u = u^q & \text{on } \partial \mathbb{R}^n_+ \times (0, T) & + \text{I.C.} \end{cases}$$
(P)

Properties of equation (P)

- A comparison argument is available.
- ► Since solutions are defined on a half space ℝⁿ₊, x_n-direction is a special one.
- ▶ Blow-up phenomenon occurs only on the boundary ∂ℝⁿ₊ under some conditions.
- Scaling invariance. $u_{\lambda}(x,t) = \lambda^{\gamma} u(\lambda x, \lambda^2 t)$.
- Energy functional.

$$E(u) = \frac{1}{2} \int_{\mathbb{R}^n_+} |\nabla u|^2 dx - \frac{1}{q+1} \int_{\partial \mathbb{R}^n_+} |u|^{q+1} dx'.$$

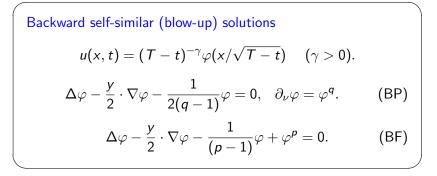
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Trivial blow-up phenomena and self-similar blow-up phenomena

$$\partial_t u = \Delta u + u^p \tag{F}$$

Trivial blow-up solutions for equation (F) $\Delta u = 0$ "

$$\partial_t u = u^p \quad \Rightarrow \quad u_{\mathsf{triv}}(t) = c_p (T-t)^{-1/(p-1)}$$



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Self-similar blow-up solutions for equation (P)

From Fila-Quittner (91), Chlebik-Fila (00), if $q \in (1, n/(n-2))$, (BP) has a unique positive bounded solution $\varphi_0(y) = \varphi_0(y_n)$:

$$arphi_0''-rac{\xi}{2}arphi_0'-rac{1}{2(q-1)}arphi_0=0 \ \ ext{in} \ \mathbb{R}_+, \quad \partial_
u arphi_0=arphi_0^q \ \ ext{on} \ \{0\}.$$

Therefore our equation (P) has a trivial blow-up solution:

$$u_{\text{triv}}(x,t) = (T-t)^{-1/2(q-1)}\varphi_0((T-t)^{-1/2}x_n).$$

Above arguments, we just constructed a special blow-up solution. How about general blow-up solutions ?

General blow-up phenomena for equation (P)

Self-similar valuables
$$x = (T - t)^{1/2}y$$
, $T - t = e^{-s}$
 $\varphi(y, s) = (T - t)^{\gamma}u(x, t)$.
blow-up time $t = T \rightarrow s = \infty$.

$$\begin{cases} \partial_{s}\varphi = \Delta \varphi - \frac{y}{2} \cdot \nabla \varphi - \frac{1}{2(q-1)}\varphi & \text{in } \mathbb{R}^{n}_{+} \times (s_{T}, \infty), \\ \partial_{\nu}\varphi = \varphi^{q} & \text{on } \partial \mathbb{R}^{n}_{+} \times (s_{T}, \infty). \end{cases}$$
(RP)

Chlebik-Fila (00) ($\gamma = 1/2(q-1), q \in (1, n/(n-2))$)

- $\varphi(y, s)$ is uniformly bounded.
- $\varphi(y, s) \to \varphi_0(y_n)$ in $C_{\text{loc}}(\overline{\mathbb{R}^n_+})$ $(s \to \infty)$. $\varphi_0(y_n)$ is a solution of (BP), that is the steady state of (RP).

Trivial blow-up profile for equation (P)

Asymptotic behavior near blow-up time

$$u(x,t) = (T-t)^{-1/2(q-1)}\varphi((T-t)^{-1/2}x, \log(T-t))$$

$$\sim (T-t)^{-1/2(q-1)}\varphi_0((T-t)^{-1/2}x_n) = u_{triv}(x,t).$$

Blow-up profile

$$U(x) = \lim_{t \to T} u(x, t) \sim \lim_{t \to T} u_{triv}(x, t)$$

= $c_q x_n^{-1/(q-1)} = c_q (\cos \theta)^{-1/(q-1)} r^{-1/(q-1)}$
=: $U_{triv}(x)$.

Polar coordinate

$$x_n = r \cos \theta$$

$$|x'| = r \sin \theta$$

$$|x'|$$

v

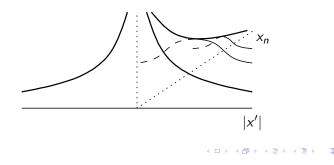
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Goal of this talk.

- What conditions on the initial data assure a single point blow-up ?
- When a single point blow-up occurs, how is the singularity of the blow-up profile on ∂ℝⁿ₊ (along |x'|-axis) ?

Single pint blow-up.



Main results

Assumption

$$x' \cdot \nabla' u_0(x) \leq 0, \quad \partial_n u_0(x) \leq 0.$$
 (A)

Theorem 1 (single point blow-up)

Let $u_0(x)$ be $u_0 \in BC^1(\overline{\mathbb{R}^n_+})$, x_n -axial symmetric and satisfy (A). Then if u(x, t) blows up in a finite time, the blow-up occurs only on the boundary.

Theorem 2 (spacial singularity of blow-up profile)

Let $u_0(x)$ be as in Theorem 1. Then if u(x, t) blows up in a finite time, u(x, t) has a blow-up profile $U(|x'|, x_n)$ and satisfies

$$c_1\left(\frac{|x'|^2}{|\log |x'||}\right)^{-1/2(q-1)} \le U(|x'|,0) \le c_2\left(\frac{|x'|^2}{|\log |x'||}\right)^{-1/2(q-1)}$$

For equation (F), Herrero-Velazquez (92-93) studied singularities of blow-up profile. We apply their strategy.

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Strategy by Herrero-Velazquez (92-93)

- 1) Dynamical system on Hilbert space $L^2_{\rho}(\mathbb{R}^n_+)$.
- 2) Improvement of convergence range.
- 3) Singularity of blow-up profile.
- 4) Selection of eigen-mode.

Our main analysis is a large time behavior of $\varphi(y, s)$.

 $\varphi(y,s) = \varphi_0(y_n) + (2'nd approximation) + (h.o.t)$

- ♦ 1'st approximation $\varphi_0(y_n) \rightarrow$ singularity of x_n -axis.
- \diamond 2'nd approximation \rightarrow singularity of |x'|-axis.

$$\begin{cases} \text{Set } v(y,s) = \varphi(y,s) - \varphi_0(y_n).\\ \partial_s v = Av, \quad \partial_\nu v = q\varphi_0^{q-1}v + f(v). \qquad (Dy) \end{cases}$$
$$Av = \Delta v - \frac{y}{2} \cdot \nabla v - \frac{1}{2(q-1)}v,\\ D(A) = \{v \in H_\rho^2(\mathbb{R}^n_+); \partial_\nu v = q\varphi_0^{q-1}v \text{ on } \partial\mathbb{R}^n_+\}.\\ \diamond \text{ Weighted Sobolev space (weight function } \rho(y) = e^{-|y|^2/4})\\ L_\rho^p(\mathbb{R}^n_+) = \left\{v \in L_{\text{loc}}^1(\mathbb{R}^n_+); \int_{\mathbb{R}^n_+} |v|^p \rho \, dy < \infty\right\},\\ H_\rho^k(\mathbb{R}^n_+) = \{v \in L_\rho^2(\mathbb{R}^n_+); D^\alpha v \in L_\rho^2(\mathbb{R}^n_+), \ \forall |\alpha| \le k\}. \end{cases}$$

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Analysis of linearized equations

Let
$$K = q\varphi_0(0)^{q-1} > 0.$$

 $\partial_s v = \Delta v - \frac{y}{2} \cdot \nabla v$ in \mathbb{R}^n_+ , $\partial_\nu v = Kv$ on $\partial \mathbb{R}^n_+$. (L1)

Set $w(y, s) = e^{\mu_K s} v(y, s) / b_K(y_n)$, where $\mu_K < 0$, $b_K(y_n)$ are the first eigenvalue and eigenfunction of

$$-\left(b'' + \frac{\xi}{2}b'\right) = \mu b \text{ in } \mathbb{R}_+, \quad \partial_\nu b = Kb \text{ on } \{0\}.$$
$$\partial_s w = \Delta w - \frac{y}{2} \cdot \nabla w + \left(\frac{2b'_K(y_n)}{b_K(y_n)}\right) \partial_n w, \quad \partial_\nu w = 0.$$
(L2)

A goal is to derive global heat kernel estimates of (L2)

$$\Gamma_{\mathcal{K}}(y,\xi,s) = \Gamma_{\mathbb{R}^{n-1}}(y',\xi',s)\gamma_{\mathcal{K}}(y_n,\xi_n,s).$$

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 $\gamma_{\mathcal{K}}(\xi,\xi',s)$ is a heat kernel of

$$\partial_s W = W_{\xi\xi} - \frac{\xi}{2} W_{\xi} + \left(\frac{2b'_{\mathcal{K}}(\xi)}{b_{\mathcal{K}}(\xi)}\right) W_{\xi}, \quad \partial_{\nu} W = 0.$$
(L3)

In the original valuable, $U(x,t)=W((1-t)^{-1/2}x,-\log(1-t)).$

$$\partial_s U = U_{xx} + (1-t)^{-1/2} \left(\frac{2b'_K}{b_K}\right) U_x, \quad \partial_\nu U = 0$$

Proposition. Let $\gamma_{\mathcal{K}}(\xi, \xi', s)$ be the heat kernel of (L3). Then there exits $\theta \in (0, 1)$ such that for $\xi \in (0, (\theta \xi' - 2)e^{s/2})$, $s \ge 1$

$$\gamma_{\mathcal{K}}(\xi,\xi',s) \leq c\gamma_{\mathcal{N}}(\xi, heta\xi',s),$$

where γ_N is the heat kernel of $\partial_s W = W_{\xi\xi} - \frac{\xi}{2}W_{\xi}$, $\partial_{\nu} W = 0$.