

Jeffery-Hamel's flow in the plane

Teppei Kobayashi
Meiji University

The 8th Japanese-German International Workshop on
Mathematical Fluid Dynamics

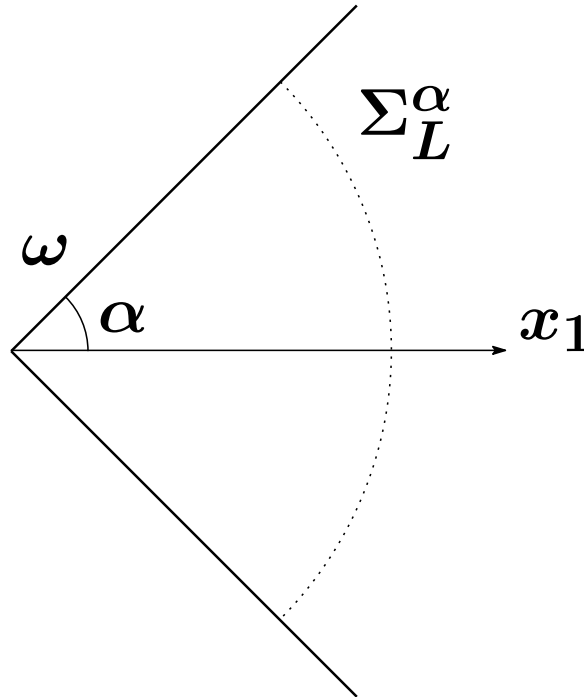
June 18, 2013 at Waseda University

1. Jeffery-Hamel's flow

$$\alpha \in (0, \pi).$$

$$\omega = \{(r, \theta); r > 0, -\alpha < \theta < \alpha\} \subset \mathbb{R}^2.$$

$$\Sigma_L^\alpha = \{(r, \theta); r = L, -\alpha < \theta < \alpha\} \quad (\forall L > 0).$$



$$\alpha = \frac{1}{2}\pi \quad \Rightarrow \quad \omega = \text{the half space.}$$

Navier-Stokes equations

$$-\nu \Delta u + (u \cdot \nabla)u + \nabla p = 0 \quad \text{in } \omega, \quad (1)$$

$$\operatorname{div} u = 0 \quad \text{in } \omega, \quad (2)$$

$$u = 0 \quad \text{on } \partial\omega, \quad (3)$$

$$u \rightarrow 0 \quad \text{as } r \rightarrow \infty \quad (4)$$

with the flux condition

$$\int_{\Sigma_L^\alpha} u \cdot e_r d\sigma = \gamma. \quad (5)$$

u : velocity

p : pressure

ν : kinematic viscosity

γ : flux constant

$e_r = (\cos \theta, \sin \theta)$

Jeffery-Hamel's flow is found in the form

$$u = \frac{\gamma g(\theta)}{r} e_r, \quad e_r = (\cos \theta, \sin \theta). \quad (6)$$

$g = g(\theta)$: unknown scalar function

ODE.

$$g'' + 4g + \frac{\gamma}{\nu} g^2 = \frac{A}{\nu\gamma} \quad \text{on} \quad (-\alpha, \alpha), \quad (7)$$

$$g(\pm\alpha) = 0, \quad (8)$$

$$\int_{-\alpha}^{\alpha} g(\theta) d\theta = 1, \quad (9)$$

$$g(\theta) = g(-\theta). \quad (10)$$

A : arbitrary constant

2. Known result (L. Rosenhead, 1940)

$$R = \frac{|\gamma|}{\nu} : \text{Reynolds number}$$

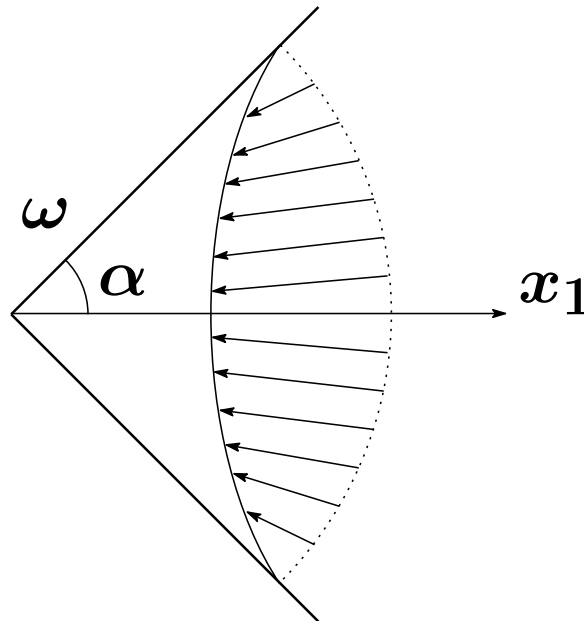
$$\gamma < 0, \alpha \in (0, \frac{\pi}{2}).$$

For $\forall R > 0$,

\Downarrow

$$\gamma g(\theta) < 0$$

$\gamma g(\theta)$ is symmetric.



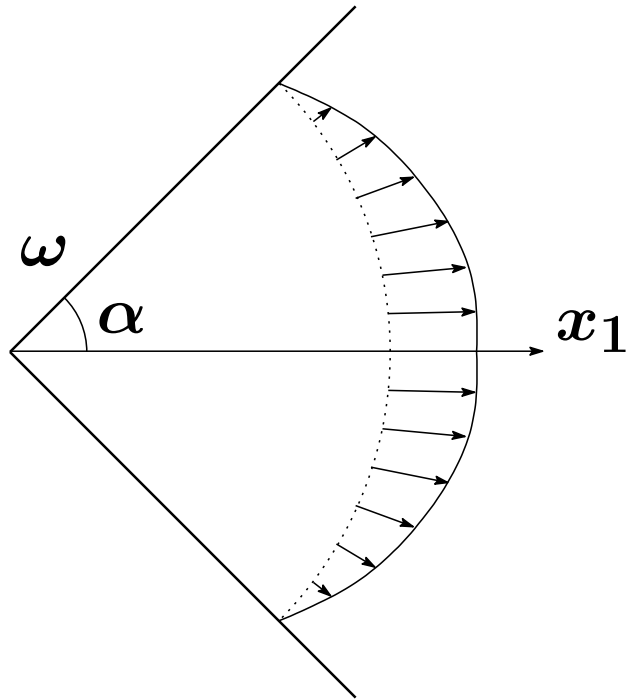
$$\gamma > 0, \alpha \in (0, \frac{\pi}{2}).$$

$\exists R_{\max}$ such that $R < R_{\max}$



$$\gamma g(\theta) > 0$$

$\gamma g(\theta)$ is symmetric.



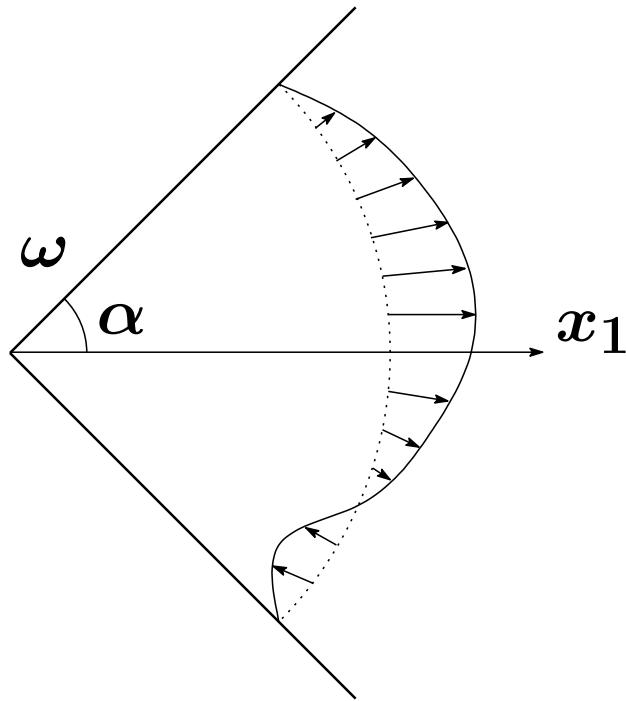
$$\gamma > 0, \alpha \in (0, \frac{\pi}{2}).$$

$R > R_{\max}$: unsteady flow



$\gamma g(\theta)$ is not always positive.

$\gamma g(\theta)$ is not symmetric.



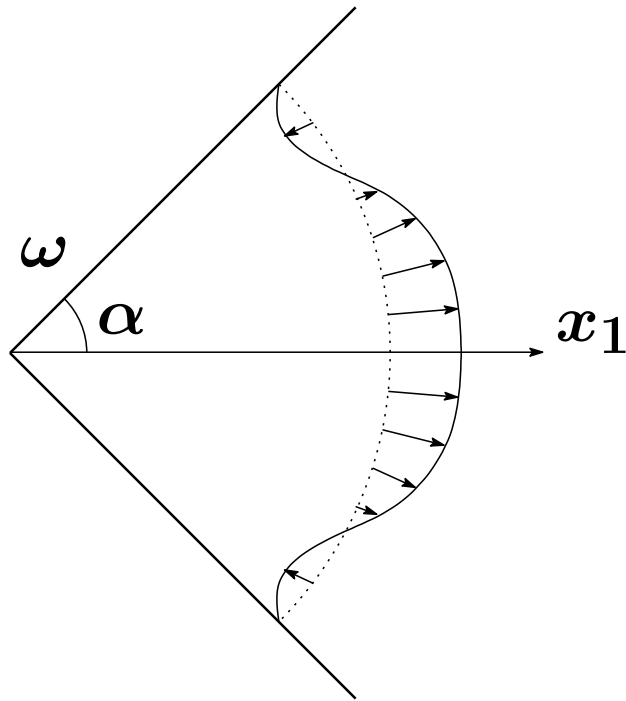
$\gamma > 0, \alpha \in (0, \frac{\pi}{2})$.

$R > R_{\max}$: unsteady flow



$\gamma g(\theta)$ is not always positive.

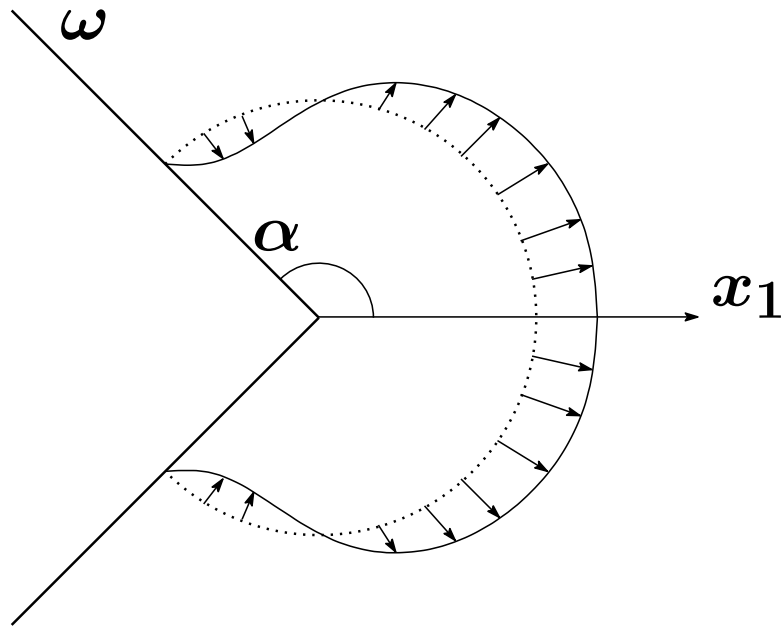
$\gamma g(\theta)$ is symmetric.



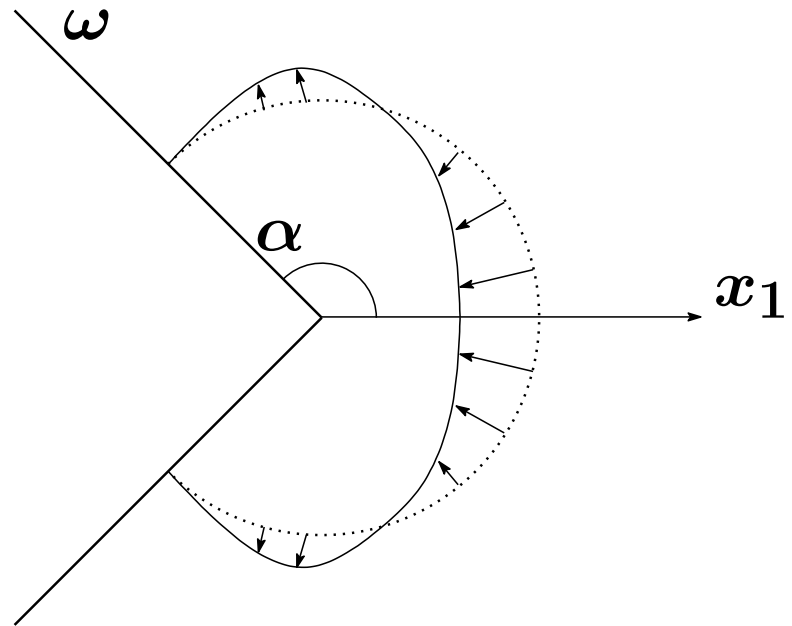
$$\alpha \in \left[\frac{\pi}{2}, \pi\right).$$

For $\forall R > 0, \theta \in (-\alpha, \alpha),$

$$\gamma > 0$$



$$\gamma < 0$$



Proposition (G. P. Galdi, M. Padula, V. A. Solonnikov(1996))

Let $\alpha = \frac{1}{2}\pi$ and $|\gamma| < \frac{\nu}{36}$.

\Downarrow

\exists solution of (N-S) in the half space of the form

$$u = \frac{\gamma g(\theta)}{r} e_r, \quad p = \frac{2\nu\gamma g(\theta) + C_1}{r^2} + C_2 \quad (C_1, C_2 \in \mathbb{R})$$

with

$$g \in C_0^S[-\frac{1}{2}\pi, \frac{1}{2}\pi] \cap C^\infty(-\frac{1}{2}\pi, \frac{1}{2}\pi), \quad \int_{-\frac{1}{2}\pi}^{\frac{1}{2}\pi} g(\theta) d\theta = 1,$$

$$|g(\theta)| \leq 6,$$

$$C_0^S[-\frac{1}{2}\pi, \frac{1}{2}\pi] = \{g \in C[-\frac{1}{2}\pi, \frac{1}{2}\pi] \mid g(\theta) = g(-\theta), g(\pm\frac{1}{2}\pi) = 0\}.$$

Moreover the function g is unique in the above class.

3. Today's result

Theorem

Let $\alpha \in (0, \pi) \setminus \{\frac{1}{4}\pi, \alpha_0(2\alpha_0 = \tan 2\alpha_0), \frac{3}{4}\pi\}$ and $|\gamma| < \gamma_0$.

\Downarrow

\exists solution of (N-S) in ω of the form

$$u = \frac{\gamma g(\theta)}{r} e_r, \quad p = \frac{2\nu\gamma g(\theta) + C_1}{r^2} + C_2 \quad (C_1, C_2 \in \mathbb{R})$$

with

$$g \in C_0^S[-\alpha, \alpha] \cap C^\infty(-\alpha, \alpha), \quad \int_{-\alpha}^{\alpha} g(\theta) d\theta = 1,$$

$$|g(\theta)| \leq \frac{2(1 - \cos 2\alpha)}{|\tan 2\alpha - 2\alpha| |\cos 2\alpha|} (:= M).$$

Moreover the function g is unique in the above class.

$$\gamma_0 = \frac{\nu\pi}{24} \quad (\alpha = \frac{1}{2}\pi).$$

$$\gamma_0 = \begin{cases} \gamma_1 & (\alpha \in (0, \frac{1}{4}\pi)), \\ \gamma_2 & (\alpha \in (\frac{1}{4}\pi, \frac{3}{4}\pi)), \\ \gamma_3 & (\alpha \in (\frac{3}{4}\pi, \pi)). \end{cases}$$

$$\gamma_1 := \frac{\nu(\tan 2\alpha - 2\alpha) \cos 2\alpha}{(1 - \cos 2\alpha + \sin 2\alpha \tan 2\alpha) \left(\frac{2\alpha(1 - \cos 2\alpha)}{(\tan 2\alpha - 2\alpha) \cos 2\alpha} + 1 \right) (1 - \cos 2\alpha)},$$

$$\gamma_2 := \frac{\nu|2\alpha - \tan 2\alpha| \cos 2\alpha}{\left(1 - \cos 2\alpha + \frac{2 - \sin 2\alpha}{|\cos 2\alpha|} \right) \left(\frac{2\alpha(1 - \cos 2\alpha)}{|2\alpha - \tan 2\alpha| \cos 2\alpha} + 1 \right) (1 - \cos 2\alpha)},$$

$$\gamma_3 := \frac{\nu(2\alpha - \tan 2\alpha) \cos 2\alpha}{\left(3 + \cos 2\alpha + \frac{4 + \sin 2\alpha}{\cos 2\alpha} \right) \left(\frac{2\alpha(1 - \cos 2\alpha)}{(2\alpha - \tan 2\alpha) \cos 2\alpha} + 1 \right) (1 - \cos 2\alpha)}.$$

5. The linear problem

For any $\alpha \in (0, \pi) \setminus \{\frac{1}{4}\pi, \frac{3}{4}\pi\}$,

$$h'' + 4h = b(\theta) \quad \text{on} \quad (-\alpha, \alpha) \quad (11)$$

with

$$h(\theta) = h(-\theta), \quad (12)$$

$$h(\pm\alpha) = 0, \quad (13)$$

where $b \in C^S[-\alpha, \alpha] = \{g \in C[-\frac{1}{2}\pi, \frac{1}{2}\pi] \mid g(\theta) = g(-\theta)\}$.

The function

$$h(\theta) = \frac{1}{2} \int_{\theta}^{\alpha} \sin 2(s - \theta) b(s) ds \\ - \frac{1}{2 \cos 2\alpha} \int_0^{\alpha} \cos 2sb(s) ds \cdot \sin 2(\alpha - \theta)$$

is a unique solution of (11) with (12) and (13).

$$\alpha = \frac{1}{4}\pi,$$

$$h'' + 4h = b(\theta) \quad \text{on} \quad (0, \frac{1}{4}\pi) \quad (14)$$

with

$$h(\frac{1}{4}\pi) = 0, \quad (15)$$

$$h'(0) = 0, \quad (16)$$

where $b \in C^S[-\alpha, \alpha]$.

The function

$$h(\theta) = \frac{1}{2} \int_{\theta}^{\alpha} \sin 2(s - \theta) b(s) ds + C_1 \sin 2\theta + C_2 \cos 2\theta$$

is a general solution of (14).

$$\begin{pmatrix} 0 & 1 \\ 0 & 2 \end{pmatrix} \begin{pmatrix} C_1 \\ C_2 \end{pmatrix} = \begin{pmatrix} 0 \\ * \end{pmatrix}. \quad \text{rank}(A) = 1, \quad \text{rank}(A, f) = 2.$$

5. The formulation

We define an operator \mathcal{L} such as

$$\mathcal{L}[b] = h \quad (b \in C^S[-\alpha, \alpha]).$$

$$g'' + 4g + \frac{\gamma}{\nu}g^2 = \frac{A}{\nu\gamma}$$

$$\mathcal{L}[g'' + 4g + \frac{\gamma}{\nu}g^2] = \frac{A}{\nu\gamma}$$

$$g = \frac{A}{\nu\gamma}\mathcal{L}[1] - \frac{\gamma}{\nu}\mathcal{L}[g^2]$$

$$\int_{-\alpha}^{\alpha} g d\theta = \frac{A}{\nu\gamma} \int_{-\alpha}^{\alpha} \mathcal{L}[1] d\theta - \int_{-\alpha}^{\alpha} \frac{\gamma}{\nu} \mathcal{L}[g^2] d\theta$$

$$1 = \frac{A}{\nu\gamma} \left(\frac{1}{2}\alpha - \frac{1}{4} \tan 2\alpha \right) - \int_{-\alpha}^{\alpha} \frac{\gamma}{\nu} \mathcal{L}[g^2] d\theta$$

$$A = \frac{\nu\gamma}{\frac{1}{2}\alpha - \frac{1}{4} \tan 2\alpha} \left(1 + \frac{\gamma}{\nu} \int_{-\alpha}^{\alpha} \mathcal{L}[g^2](\theta) d\theta \right)$$

$$g = \frac{1}{\frac{1}{2}\alpha - \frac{1}{4} \tan 2\alpha} \left(1 + \frac{\gamma}{\nu} \int_{-\alpha}^{\alpha} \mathcal{L}[g^2](\theta) d\theta \right) \mathcal{L}[1] - \frac{\gamma}{\nu} \mathcal{L}[g^2]$$

$(g \in C_0^S[-\alpha, \alpha]).$

For any $\alpha \in (0, \pi) \setminus \{\frac{1}{4}\pi, \alpha_0(2\alpha_0 = \tan 2\alpha_0), \frac{3}{4}\pi\}$, we set

$$\mathcal{J}[g] := \frac{1}{\frac{1}{2}\alpha - \frac{1}{4}\tan 2\alpha} \left(1 + \frac{\gamma}{\nu} \int_{-\alpha}^{\alpha} \mathcal{L}[g^2](\theta) d\theta \right) \mathcal{L}[1] - \frac{\gamma}{\nu} \mathcal{L}[g^2]$$

$(g \in C_0^S[-\alpha, \alpha]).$

Lemma

Let $|\gamma| < \gamma_0$.

\Downarrow

Then \mathcal{J} is a contraction from $B(0, M)$ to $B(0, M)$,

where $B(0, M) = \{\varphi \in C_0^S[-\alpha, \alpha]; \|\varphi\|_C \leq M\}$,

$$M = \frac{2(1 - \cos 2\alpha)}{|\tan 2\alpha - 2\alpha| |\cos 2\alpha|}.$$