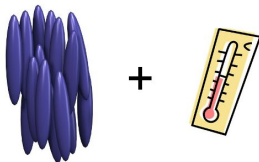


# A thermodynamically consistent extension of the simplified Ericksen-Leslie Model

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jointly with M. Hieber and J. Prüss



9th Japanese-German International Workshop on  
Mathematical Fluid Dynamics

Waseda University  
Nov 5, 2013

# 1. The *old* Model

- ▶  $\Omega$ : bounded  $C^2$ -domain.

$$\left( \text{blue scribble} \right) \left\{ \begin{array}{ll} \partial_t u + (u \cdot \nabla)u - \mu \Delta u + \nabla \pi = -\lambda \operatorname{div}([\nabla d]^T \nabla d) & (0, T) \times \Omega, \\ \partial_t d + (u \cdot \nabla)d = \gamma(\Delta d + |\nabla d|^2 d) & (0, T) \times \Omega, \\ |d| = 1 & (0, T) \times \Omega, \\ \operatorname{div} u = 0 & (0, T) \times \Omega, \\ (u, \partial_\nu d) = (0, 0) & (0, T) \times \partial\Omega, \\ (u, d)|_{t=0} = (u_0, d_0) & \Omega. \end{array} \right.$$

$u : (0, \infty) \times \Omega \rightarrow \mathbb{R}^n$ : velocity,

$\pi : (0, \infty) \times \Omega \rightarrow \mathbb{R}$ : pressure,

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How to realize this in an extended system  ?

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$$\text{where } D := \frac{1}{2}([\nabla u]^T + \nabla u), \quad \mathbb{S} := 2\mu D - \lambda[\nabla d]^T \nabla d.$$

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## 4. Internal Energy of new System



### Total energy

$$E = E_{kin} + E_{int} = \int_{\Omega} \frac{1}{2} |u|^2 + e \, dx,$$

where  $e = e(\vartheta, \frac{|\nabla d|^2}{2})$ : internal energy.

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### Internal relations

$$\left\{ \begin{array}{l} e = \psi + \vartheta \eta, \\ \eta = -\partial_1 \psi, \quad \kappa_0 = \partial_1 e, \quad \lambda = \partial_2 \psi, \\ \psi'' > 0. \end{array} \right.$$

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### Connection to -d) and - $\vartheta$ )

$$\begin{aligned} (\partial_t + u \cdot \nabla) e &= \partial_1 e (\partial_t + u \cdot \nabla) \vartheta + \partial_2 e (\partial_t + u \cdot \nabla) (|\nabla d|^2 / 2). \\ &= \kappa_0 (\partial_t + u \cdot \nabla) \vartheta + (\lambda - \vartheta \partial_1 \lambda) (\partial_t + u \cdot \nabla) (|\nabla d|^2 / 2). \end{aligned}$$

## 5. Clever RHS for ( - $\vartheta$ ) leads to §1

### Idea

- ▶ Choose right-hand side of ( - $\vartheta$ ) such that

$$\partial_t \int_{\Omega} e \, dx = -\partial_t \int_{\Omega} \frac{1}{2} |u|^2 \, dx.$$



## 5. Clever RHS for $(\vartheta - \vartheta)$ leads to $\xi_1$

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### Suitable right-hand side

$$\begin{aligned} (\vartheta - \vartheta) \quad & \kappa_0(\partial_t \vartheta + (u \cdot \nabla) \vartheta) + \operatorname{div} q = 2\mu |D|^2 - \partial_1 \lambda \vartheta ([\nabla d]^T \nabla d) : \nabla u \\ & - (\lambda - \vartheta \partial_1 \lambda) \nabla d : \nabla [\gamma (\vartheta \operatorname{div} (\frac{\lambda}{\vartheta} \nabla d) + \lambda |\nabla d|^2 d)] \end{aligned}$$

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- ▶ Boundary condition  $q \cdot n = 0$ .

## 6. Clausius-Duhem Assumption leads to §2


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(-CL) 
$$\left\{ \begin{array}{l} \int_{\Omega} (\partial_t + u \cdot \nabla) \eta + \operatorname{div} \left( \frac{q}{\vartheta} \right) dx \geq 0, \\ \text{with equality } \textit{iff} \text{ the system is in an equilibrium state.} \end{array} \right.$$

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- ▶ Local (strong) well-posedness,
- ▶ Solution is “as regular” as physical coefficients (Angenent’s trick).

## Asymptotic behavior

- ▶ Exponential stability for solutions with initial data close to equilibria.

# 7. Consequences

## Modeling

- ▶  $\mathcal{M}$  is *a priori* extension of  $\mathcal{M}_0$ .
- ▶  $\mathcal{M}$  is *a priori* thermodynamically consistent.
- ▶ Lyapunov functional “for free” ( $\rightarrow$  asymptotics).

As parabolic quasilinear equation one obtains

## Dynamics

- ▶ Local (strong) well-posedness,
- ▶ Solution is “as regular” as physical coefficients (Angenent’s trick).

## Asymptotic behavior

- ▶ Exponential stability for solutions with initial data close to equilibria.
- ▶ Exponential stability for solutions which are eventually bounded on their maximal existence interval.