



JOHANNES GUTENBERG
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Numerical tests on some viscoelastic flows

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An overview of different models

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Navier Stokes

$$\begin{cases} Re(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u}) = -\nabla p + \nabla \cdot (2\alpha \mathbf{D}(\mathbf{u}) + \boldsymbol{\sigma}) \\ \nabla \cdot \mathbf{u} = 0 \end{cases}$$

\mathbf{u} : velocity,

p : pressure,

Re : Reynolds number,

α : ratio of Newtonian viscosity,

$\boldsymbol{\sigma}$: elastic viscosity ,

$$\mathbf{D}(\mathbf{u}) = \frac{\nabla \mathbf{u} + \nabla \mathbf{u}^T}{2}.$$

Oldroyd-B model

$$\frac{\partial \boldsymbol{\tau}}{\partial t} + (\mathbf{u} \cdot \nabla) \boldsymbol{\tau} - \nabla \mathbf{u} \cdot \boldsymbol{\tau} - \boldsymbol{\tau} \cdot (\nabla \mathbf{u})^T = \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau})$$

where $\boldsymbol{\sigma} = \frac{1-\alpha}{We} (\boldsymbol{\tau} - \mathbf{I})$.

$\boldsymbol{\tau}$: conformation tensor, symmetric positive definite

We : Weissenberg number, relaxation time over characteristic time

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where $\boldsymbol{\sigma} = \frac{1}{We} (tr \boldsymbol{\tau}) \boldsymbol{\tau}$.

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Pertlerin model

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model	global weak solution	numerics
Oldroyd-B	no	no
Rotational	yes	yes
Pertlerin	yes	yes

some stability results

- $\psi = \text{sqrt}(\tau)$ Square root(SR)
- $\psi = \text{log}(\tau)$ Logarithm conformation representation(LCR)



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If τ is symmetric positive definite, matrix $\nabla \mathbf{u}$ can be decomposed as following

$$\nabla \mathbf{u} = \mathbf{B} + \mathbf{\Omega} + \mathbf{N}\tau^{-1}.$$

where \mathbf{N} , $\mathbf{\Omega}$ anti-symmetric, \mathbf{B} symmetric and commutes with τ .

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$$R^T \tau R = \text{diag}(\lambda_1, \lambda_2)$$

$$\begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} = R^T (\nabla \mathbf{u}) R.$$

$$\mathbf{N} = R \begin{pmatrix} 0 & n \\ -n & 0 \end{pmatrix} R^T, \mathbf{B} = R \begin{pmatrix} m_{11} & 0 \\ 0 & m_{22} \end{pmatrix} R^T, \mathbf{\Omega} = R \begin{pmatrix} 0 & \omega \\ -\omega & 0 \end{pmatrix} R^T$$

with $n = (m_{12} + m_{21})/(\lambda_2^{-1} - \lambda_1^{-1})$, $\omega = (\lambda_2 m_{12} + \lambda_1 m_{21})/(\lambda_2 - \lambda_1)$.

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$$\nabla \mathbf{u} \cdot \tau + \tau \cdot (\nabla \mathbf{u})^T = \mathbf{\Omega}\tau - \tau\mathbf{\Omega} + 2\mathbf{B}\tau$$

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$$\frac{\partial \boldsymbol{\tau}}{\partial t} + \mathbf{u} \cdot \boldsymbol{\tau} = \boldsymbol{\Omega} \boldsymbol{\tau} - \boldsymbol{\tau} \boldsymbol{\Omega} + 2\mathbf{B} \boldsymbol{\tau} + \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau})$$

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Note:

If \mathbf{Y} is skew-symmetric, $e^{\mathbf{Y}}$ is orthogonal,

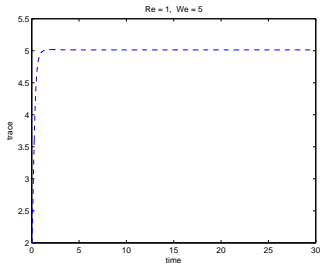
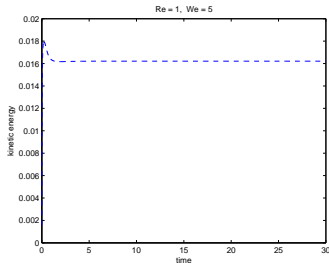
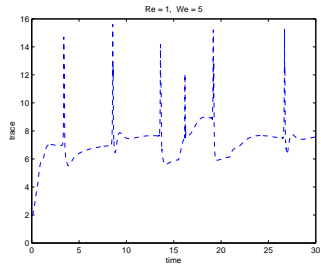
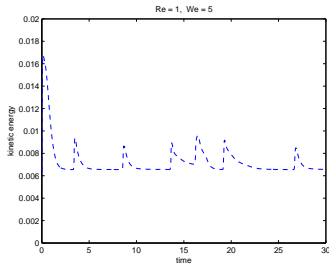
If \mathbf{Y} is invertible, then $e^{\mathbf{Y} \mathbf{X} \mathbf{Y}^{-1}} = \mathbf{Y} e^{\mathbf{X}} \mathbf{Y}^{-1},$

If $\mathbf{X} \mathbf{Y} = \mathbf{Y} \mathbf{X},$ then $e^{\mathbf{X}} e^{\mathbf{Y}} = e^{\mathbf{X} + \mathbf{Y}}.$

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- SR stable for relative high Weissenberg number,
LCR stable for arbitrary high Weissenberg number



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EOC for Oldroyd-B model with LCR

$$e(\phi(N)) = \|\phi(N) - \phi(256)\|_2$$



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1. for low We

Table : We=0.5

t	$e(\psi_{11}^{64})$	$e(\psi_{11}^{128})$	EOC	$e(\tau_{11}^{64})$	$e(\tau_{11}^{128})$	EOC
1	0.0722	0.0310	1.2225	0.6398	0.3212	0.9942
2	0.0864	0.0381	1.1808	1.0305	0.5773	0.8358
4	0.0864	0.0381	1.1664	1.3045	0.7836	0.7354
8	0.0928	0.0414	1.1657	1.3245	0.8001	0.7273

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2. for high We

Table : We=2

t	$e(\psi_{11}^{64})$	$e(\psi_{11}^{128})$	EOC	$e(\tau_{11}^{64})$	$e(\tau_{11}^{128})$	EOC
1	0.0989	0.0428	1.2095	1.2414	0.6129	1.0182
2	0.1425	0.0630	1.1767	6.8478	4.1627	0.7181
4	0.1961	0.0882	1.1525	20.5848	15.0224	0.4545

characteristic method



characteristic method

material derivative

$$\frac{\partial f}{\partial t} + \mathbf{u} \cdot \nabla f \rightarrow \frac{Df}{Dt} = \frac{f^n - f^{n-1} \circ X_1(u^{n-1}, \Delta t)}{\Delta t} + O(\Delta t)$$

$$X = X_1 + u^{n-1} \Delta t$$

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- finite difference
- finite element

Part I.

$$\mathbf{u}_t + \mathbf{u} \cdot \nabla \mathbf{u} = \frac{\alpha}{Re} \Delta \mathbf{u} + \frac{1-\alpha}{Re} \frac{1}{We} \nabla \cdot \boldsymbol{\tau}$$

$$\mathbf{u}_t = -\nabla p \quad (\nabla \cdot \mathbf{u} = 0)$$

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$$U^* - \Delta t \frac{\alpha}{Re} (U_{xx}^* + U_{yy}^*) = U^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t) + \Delta t \frac{1-\alpha}{Re We} (\tau_{11x}^{k-1} + \tau_{12y}^{k-1})$$
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$$\frac{U^k - U^*}{\Delta t} = -(P^k)_x$$
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Part II.

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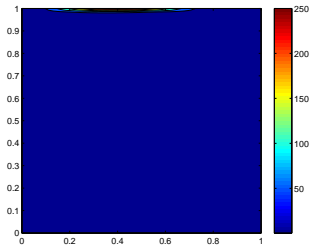
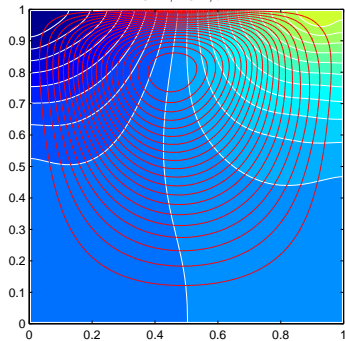
Part II.

$$\frac{\psi^k - \psi^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)}{\Delta t} = \boldsymbol{\Omega}^{k-1} \psi^{k-1} - \psi^{k-1} \boldsymbol{\Omega}^{k-1} + 2\mathbf{B}^{k-1} + \frac{e^{-\psi^{k-1}} - \mathbf{I}}{We}$$

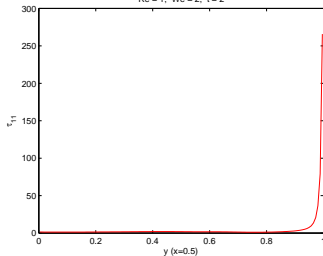
Weakly coupled, iteration over $\mathbf{u}, p, \boldsymbol{\tau}$.



Re = 1, We=2, t = 2

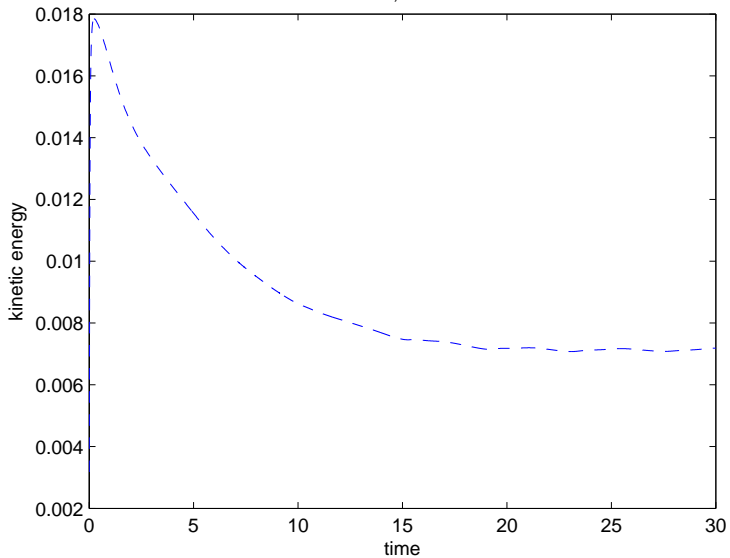


Re = 1, We = 2, t = 2





Re = 1, We = 5



Weak formula:

$$\left\{ \begin{array}{l} Re \frac{\mathbf{u}^k - \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)}{\Delta t} + \nabla p^k - 2\alpha \nabla \cdot \mathbf{D}(\mathbf{u}^k) = \frac{1-\alpha}{We} \nabla \cdot (\boldsymbol{\tau}^{k-1} - \mathbf{I}) \\ \nabla \cdot \mathbf{u}^k = 0 \\ \frac{\boldsymbol{\tau}^k - \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^k, \Delta t)}{\Delta t} = \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau}^k) \end{array} \right.$$

Weak formula:

$$\left\{ \begin{array}{l} Re \frac{\mathbf{u}^k - \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)}{\Delta t} + \nabla p^k - 2\alpha \nabla \cdot \mathbf{D}(\mathbf{u}^k) = \frac{1-\alpha}{We} \nabla \cdot (\boldsymbol{\tau}^{k-1} - \mathbf{I}) \\ \nabla \cdot \mathbf{u}^k = 0 \\ \frac{\boldsymbol{\tau}^k - \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^k, \Delta t)}{\Delta t} = \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau}^k) \end{array} \right.$$

$$\begin{aligned} & \frac{Re}{\Delta t} (v, u) + 2\alpha (\mathbf{D}(v), \mathbf{D}(u)) - (\nabla \cdot v, p) - (\nabla \cdot u, q) \\ &= (v, \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) + \frac{1-\alpha}{We} (v, \nabla \cdot \boldsymbol{\tau}^{k-1}) \\ & \left(\frac{1}{\Delta t} + \frac{1}{We} \right) (\phi, \boldsymbol{\tau}^k) \\ &= (\phi, \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} \mathbf{I}) + (\phi, \frac{1}{\Delta t} \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) \end{aligned}$$

Weak formula:

$$\begin{cases} Re \frac{\mathbf{u}^k - \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)}{\Delta t} + \nabla p^k - 2\alpha \nabla \cdot \mathbf{D}(\mathbf{u}^k) = \frac{1-\alpha}{We} \nabla \cdot (\boldsymbol{\tau}^{k-1} - \mathbf{I}) \\ \nabla \cdot \mathbf{u}^k = 0 \\ \frac{\boldsymbol{\tau}^k - \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^k, \Delta t)}{\Delta t} = \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau}^k) \end{cases}$$

$$\begin{aligned} & \frac{Re}{\Delta t} (v, u) + 2\alpha (\mathbf{D}(v), \mathbf{D}(u)) - (\nabla \cdot v, p) - (\nabla \cdot u, q) \\ &= (v, \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) + \frac{1-\alpha}{We} (v, \nabla \cdot \boldsymbol{\tau}^{k-1}) \\ & \left(\frac{1}{\Delta t} + \frac{1}{We} \right) (\phi, \boldsymbol{\tau}^k) \\ &= (\phi, \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} \mathbf{I}) + (\phi, \frac{1}{\Delta t} \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) \end{aligned}$$

$P2/P1/(P1)^3$ elements,
 $(P1/P1 + \text{pressure stabilization})/(P1)^3$ elements.

Weak formula:


$$\begin{cases} Re \frac{\mathbf{u}^k - \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)}{\Delta t} + \nabla p^k - 2\alpha \nabla \cdot \mathbf{D}(\mathbf{u}^k) = \frac{1-\alpha}{We} \nabla \cdot (\boldsymbol{\tau}^{k-1} - \mathbf{I}) \\ \nabla \cdot \mathbf{u}^k = 0 \\ \frac{\boldsymbol{\tau}^k - \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^k, \Delta t)}{\Delta t} = \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} (\mathbf{I} - \boldsymbol{\tau}^k) \end{cases}$$

$$\begin{aligned} & \frac{Re}{\Delta t} (v, u) + 2\alpha (\mathbf{D}(v), \mathbf{D}(u)) - (\nabla \cdot v, p) - (\nabla \cdot u, q) \\ &= (v, \mathbf{u}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) + \frac{1-\alpha}{We} (v, \nabla \cdot \boldsymbol{\tau}^{k-1}) \\ & \left(\frac{1}{\Delta t} + \frac{1}{We} \right) (\phi, \boldsymbol{\tau}^k) \\ &= (\phi, \nabla \mathbf{u}^k \boldsymbol{\tau}^{k-1} + \boldsymbol{\tau}^{k-1} (\nabla \mathbf{u}^k)^T + \frac{1}{We} \mathbf{I}) + (\phi, \frac{1}{\Delta t} \boldsymbol{\tau}^{k-1} \circ X_1(\mathbf{u}^{k-1}, \Delta t)) \end{aligned}$$

$P2/P1/(P1)^3$ elements,

$(P1/P1 + \text{pressure stabilization})/(P1)^3$ elements.

$$\frac{\boldsymbol{\psi}^k - \boldsymbol{\psi}^{k-1} \circ X_1(\mathbf{u}^k, \Delta t)}{\Delta t} = \Omega^{k-1} \boldsymbol{\psi}^{k-1} + \boldsymbol{\psi}^{k-1} \Omega + \frac{1}{We} (e^{-\boldsymbol{\psi}^{k-1} (\nabla \mathbf{u}^k)^T \boldsymbol{a}} - \mathbf{I})$$



Thank you for your attention!