

Quasi-optimal initial value conditions for the Navier-Stokes equations

11th Japanese-German International Workshop

on Mathematical Fluid Dynamics,

Waseda University, Tokyo, March 10-13, 2015

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Navier–Stokes Equations (NSt)

$$\begin{aligned}u_t - \Delta u + u \cdot \nabla u + \nabla p &= f && \text{in } \Omega \times (0, T) \\ \operatorname{div} u &= 0 && \text{in } \Omega \times (0, T) \\ u &= 0 && \text{on } \partial\Omega \times (0, T) \\ u(0) &= u_0 && \text{at } t = 0\end{aligned}$$

where $\Omega \subset \mathbb{R}^3$ bounded (or exterior domain), $\partial\Omega \in C^{1,1}$,
 $0 < T \leq \infty$, $\nu = 1$

Theorem (Existence of Weak Leray-Hopf Solutions)

Let $u_0 \in L^2_\sigma(\Omega)$ and $f = \operatorname{div} F$ with $F \in L^2(0, T; L^2)$ be given.

Then there exists at least one *weak Leray-Hopf solution*

$$u \in \mathcal{LH}_T = L^\infty(0, T; L^2_\sigma(\Omega)) \cap L^2(0, T; H^1_{0,\sigma}(\Omega))$$

of (NSt) in the sense of distributions. Moreover,

$$u \in C_w^0([0, T]; L^2_\sigma(\Omega)), \text{ i.e., } u \text{ is weakly continuous in } L^2_\sigma(\Omega),$$

and u satisfies the *strong energy inequality (SEI)*, i.e.,

$$\frac{1}{2} \|u(t)\|_2^2 + \int_{t_0}^t \|\nabla u\|_2^2 d\tau \leq \frac{1}{2} \|u(t_0)\|_2^2 + \int_{t_0}^t \langle f, u \rangle d\tau$$

for a.a. $0 \leq t_0 \leq T$ (including $t_0 = 0$) and all $t \in (t_0, T)$.

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Note A weak solution satisfies

$$u \in L^s(0, T; L^q_\sigma(\Omega)) \quad \text{for all } s \geq 2, q \geq 2 : \frac{2}{s} + \frac{3}{q} = \frac{3}{2}$$

Strong Solutions

Definition A weak solution $u \in \mathcal{LH}_T$ is called a *strong solution* if additionally

$$u \in L^s(0, T; L^q_\sigma(\Omega)) \quad \text{for some } s > 2, q > 3 : \frac{2}{s} + \frac{3}{q} = 1$$

Theorem (Uniqueness) Let u be a strong solution of (NSt). Then u is unique among all weak solutions satisfying (EI) (only $t_0 = 0$ is required).

Theorem (Regularity) Let u be a strong solution of (NSt) and $f \in C_0^\infty(\bar{\Omega} \times (0, T))$, $\partial\Omega \in C^\infty$. Then $u \in C^\infty(\bar{\Omega} \times (0, T))$.

Let $\Omega \subset \mathbb{R}^3$ bounded with $\partial\Omega \in C^{1,1}$ and $u_0 \in L^2_\sigma(\Omega)$, let $2 < s, 3 < q$ and $\frac{2}{s} + \frac{3}{q} = 1$.

① The condition

$$\int_0^\infty \|e^{-\tau A_2} u_0\|_q^s d\tau < \infty$$

is *necessary and sufficient* for the existence of a unique strong solution $u \in L^s(0, T; L^q)$, $0 < T < \infty$, of (NSt).

Let $\Omega \subset \mathbb{R}^3$ bounded with $\partial\Omega \in C^{1,1}$ and $u_0 \in L^2_\sigma(\Omega)$, let $2 < s$, $3 < q$ and $\frac{2}{s} + \frac{3}{q} = 1$.

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- ② $\exists \varepsilon_* = \varepsilon_*(\Omega, q) > 0$: If

$$\int_0^T \|e^{-\tau A_2} u_0\|_q^s d\tau \leq \varepsilon_*$$

then there exists a unique strong solution u on $[0, T)$ with $u(0) = u_0$.

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Exterior domain: Komo-Fa., Analysis 33 (2013)

General domain: L^2 -theory (Sohr, Varnhorn, Fa. (2009)) or \tilde{L}^q -spaces $\tilde{L}^q = L^q \cap L^2$ (Riechwald, PhD, Darmstadt 2011)

Ideas of Proof I

(i) A weak solution u has the representation

$$u(t) = E(t) - \int_0^t A^{1/2} e^{-(t-\tau)A} (A^{-1/2} P \operatorname{div})(u \otimes u)(\tau) d\tau$$

where $E(t) = e^{-tA} u_0$ and the formal operator $A^{-1/2} P \operatorname{div}$ is bounded on each $L^q(\Omega)$, $1 < q < \infty$

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(ii) The operator \mathcal{F} for $\tilde{u} = u - E$,

$$\mathcal{F}(\tilde{u})(t) = - \int_0^t A^{1/2} e^{-(t-\tau)A} (A^{-1/2} P \operatorname{div})(\tilde{u} + E) \otimes (\tilde{u} + E)(\tau) d\tau,$$

defines a strict contraction in a closed ball of $L^s(0, T'; L^q(\Omega))$ for sufficiently small $0 < T' < \infty$ (need ε_* and T' small)

(iii) Show that $\tilde{u} = \mathcal{F}(\tilde{u})$ defines a weak solution $u = \tilde{u} + E$ of (NSt) in $\mathcal{LH}_{T'}$ with initial value u_0

Step 1: Show that $\nabla \tilde{u} \in L^2(0, T; L^2(\Omega))$:

Use Yosida approximation operators to get for $\tilde{u}_n = J_n \tilde{u}$ that

$$\left\| A^{\frac{1}{2}} \tilde{u}_n \right\|_{2,2;T} \leq c \|u\|_{q,s;T} \left(\left\| A^{\frac{1}{2}} \tilde{u}_n \right\|_{2,2;T} + \|\nabla E\|_{2,2;T} \right)$$

Choose T small for absorption \Rightarrow uniformly in $n \in \mathbb{N}$

$$\left\| A^{\frac{1}{2}} \tilde{u}_n \right\|_{2,2;T} \leq 2c \|u\|_{q,s;T} \|\nabla E\|_{2,2;T} < \infty.$$

$n \rightarrow \infty \Rightarrow$

$$A^{\frac{1}{2}} \tilde{u}, \nabla \tilde{u}, \nabla u \in L^2(0, T'; L^2(\Omega))$$

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Step 2: Show that $uu \in L^2(0, T; L^2) \Rightarrow u$ is a weak Leray-Hopf solution satisfying the energy equality on $[0, T]$ □

Besov space notation by H. Amann (2002):

$$\mathbb{B}_{q',s'}^{2/s}(\Omega) = B_{q',s'}^{2/s}(\Omega) \cap L_{\sigma}^{q'}(\Omega) = \{v \in B_{q',s'}^{2/s}(\Omega) : \operatorname{div} v = 0, N \cdot v|_{\partial\Omega} = 0\}$$

Define

$$\mathbb{B}_{q,s}^{-2/s}(\Omega) = (\mathbb{B}_{q',s'}^{2/s}(\Omega))^* = (L_{\sigma}^{q'}, \mathcal{D}(A_{q'}))^*_{1-1/s',s'} = (\mathcal{D}(A_{q'})^*, L_{\sigma}^q)_{1-1/s,s}$$

using real interpolation, duality \Rightarrow

$$\|u_0\|_{\mathbb{B}_{q,s}^{-2/s}(\Omega)} \sim \|A^{-1}u_0\|_{(L_{\sigma}^q(\Omega), \mathcal{D}(A_q))_{1-1/s,s}} \sim \left(\int_0^{\delta} \|e^{-\tau A}u_0\|_q^s d\tau \right)^{1/s}$$

for bounded Ω and with δ

\Rightarrow integrability of $\|e^{-\tau A}u_0\|_q^s$, $u_0 \in L_{\sigma}^2(\Omega)$, is crucial only near 0

Applications

Apply criteria using $\mathbb{B}_{q,s}^{-2/s}(\Omega)$ along a given weak solution at all t_0 (or at almost all t_0 (?)) and identify the local strong solution starting at t_0 with the weak solution via Serrin's Uniqueness Theorem

⇒ left-hand or right-hand local/global regularity and uniqueness, see Sohr, Varnhorn, F. (Indiana Univ. Math. J. 2007, JMFM 2008, 2012, 2014, Ann. Univ. Ferrara 2009, 2014)

Extension to Quasi-optimal Initial Values

$$\mathcal{D}(A_2^{1/4}) \subset L_\sigma^3(\Omega) \subset L_\sigma^{3,r}(\Omega) \subset \mathbb{B}_{q,s_q}^{-2/s_q}(\Omega) = \mathbb{B}_{q,s_q}^{-1+3/q}(\Omega)$$

where $\frac{2}{s_q} + \frac{3}{q} = 1$.

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For $s_q < s < \infty$ we proceed with

$$\mathbb{B}_{q,s_q}^{-1+3/q}(\Omega) \subset \mathbb{B}_{q,s}^{-1+3/q}(\Omega) \subset \mathbb{B}_{q,\infty}^{-1+3/q}(\Omega)$$

where $\frac{2}{s} + \frac{3}{q} = 1 - 2\alpha$, $0 < \alpha < \frac{1}{2}$.

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Here $\mathbb{B}_{q,s}^{-1+3/q}(\Omega) = (\mathcal{D}(A_{q'})^*, L_\sigma^q(\Omega))_{1-\alpha-1/s,s}$ and

$$\begin{aligned} \|u_0\|_{\mathbb{B}_{q,s}^{-1+3/q}} &\sim \|e^{-\tau A} u_0\|_{L_\alpha^s(L^q)} := \int_0^\infty (\tau^\alpha \|e^{-\tau A} u_0\|_q)^s d\tau < \infty \\ &\sim \int_0^\delta (\tau^\alpha \|e^{-\tau A} u_0\|_q)^s d\tau < \infty \end{aligned}$$

How far do the previous results hold when $s_q < s < \infty$?

We are looking for an $L_\alpha^s(L^q)$ -strong solution u , i.e., u is a Leray-Hopf weak solution and

$$\int_0^T (\tau^\alpha \|u(\tau)\|_q)^s d\tau < \infty$$

where $2 < s < \infty$, $3 < q < \infty$, $0 < \alpha < \frac{1}{2}$ with $\frac{2}{s} + \frac{3}{q} = 1 - 2\alpha$

Main Theorem (Giga, Hsu, F. ('14))

Let $u_0 \in L^2_\sigma(\Omega)$ and s, q, α be given.

(1) There exists an $\epsilon_* = \epsilon_*(q, s, \alpha, \Omega) > 0$: If

$$\|e^{-\tau A} u_0\|_{L^s_\alpha(0, T; L^q)} \leq \epsilon_*,$$

then (NSt) has a unique $L^s_\alpha(L^q)$ -strong solution with data u_0 on $[0, T)$.

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(2) The condition

$$\int_0^\infty (\tau^\alpha \|e^{-\tau A} u_0\|_q)^s d\tau < \infty$$

is sufficient and necessary for the existence of a unique $L^s_\alpha(L^q)$ -strong solution $u \in L^s_\alpha(0, T; L^q)$ with data u_0 .

(3) The $L^s_\alpha(L^q)$ -strong solution satisfies the **energy equality on $[0, T)$** .

Ideas of Proof I

(i) Construction of the $L_\alpha^s(L^q)$ -strong solution: Find $\tilde{u} = u - E$, $E(t) = e^{-tA}u_0$, as fixed point $\tilde{u} = \mathcal{F}(\tilde{u})$ of the operator

$$\mathcal{F}(\tilde{u})(t) = - \int_0^t A^{1/2} e^{-(t-\tau)A} (A^{-1/2} P \operatorname{div})(\tilde{u} + E) \otimes (\tilde{u} + E)(\tau) d\tau.$$

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Lemma: (Weighted Hardy-Littlewood-Sobolev inequality)

Let $0 < \lambda < 1$, $\alpha_2 \leq \alpha_1$, $1 < s_1 \leq s_2 < \infty$, $-\frac{1}{s_j} < \alpha_j < 1 - \frac{1}{s_j}$, $j = 1, 2$ and $\frac{1}{s_1} + (\lambda + \alpha_1 - \alpha_2) = 1 + \frac{1}{s_2}$. Then

$$I_\lambda f(t) = \int_{\mathbb{R}} (t - \tau)^{-\lambda} f(\tau) d\tau$$

defines a bounded operator $I_\lambda : L_{\alpha_1}^{s_1}(\mathbb{R}) \rightarrow L_{\alpha_2}^{s_2}(\mathbb{R})$.

Ideas of Proof II

Lemma: $\nabla u \in L^2(0, T; L^2)$, hence $u \in L^2(0, T; L^6)$

Proof Weighted Hardy-Littlewood-Sobolev inequality \Rightarrow

$$\left\| A^{\frac{1}{2}} J_n u \right\|_{2,2;T} \leq 2c \|u\|_{L_\alpha^s(L^q)} \|\nabla E\|_{2,2;T} < \infty. \quad \square$$

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As a direct consequence of Hölder's inequality and $\tilde{u} = \mathcal{F}(\tilde{u})$ (this argument requires $\alpha > 0$ (!))

$$\|\tilde{u}(t)\|_2 \leq c \|u\|_{L_\alpha^s(0,t;L^q)} \|\nabla u\|_{L^2(L^2)}$$

$\Rightarrow u \in L^\infty(0, T; L^2)$ and $\|\tilde{u}(t)\|_2 \rightarrow 0$ as $t \rightarrow 0$

$\Rightarrow u(t) \rightarrow u_0$ in $L^2(\Omega)$ as $t \rightarrow 0$

$\Rightarrow u$ is a weak solution with $u(0) = u_0$

Ideas of Proof III

Lemma: $u \in L^4_{\alpha/(2+8\alpha)}(0, T; L^4(\Omega))$.

Proof With $\beta = \frac{1}{2+8\alpha}$, $q_1, s_1 \dots, \frac{2}{s_1} + \frac{3}{q_1} = \frac{3}{2}$, and an interpolation inequality

$$\int_0^T \tau^{4\alpha\beta} \|u\|_4^4 d\tau \leq \|u\|_{L^s_\alpha(L^q)}^{4\beta} \|u\|_{L^{s_1}(L^{q_1})}^{4(1-\beta)} < \infty. \quad \square$$

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$\Rightarrow u \in L^4(\varepsilon, T; L^4(\Omega)) \Rightarrow u \otimes u \in L^2(\varepsilon, T; L^2)$ can be considered as right-hand side of a Stokes system on (ε, T)

$\Rightarrow u$ satisfies the energy equality on (ε, T) :

$$\frac{1}{2} \|u(t)\|_2^2 + \int_\varepsilon^t \|\nabla u\|_2^2 d\tau = \frac{1}{2} \|u(\varepsilon)\|_2^2$$

Since $u(\varepsilon) \rightarrow u_0$ in $L^2(\Omega)$ as $\varepsilon \rightarrow 0$,

u satisfies the energy equality on $[0, T)$.

Are $L^s_\alpha(L^q)$ -strong solutions unique?

The proof of Serrin's uniqueness theorem

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Uniqueness Theorem: Let $2 < s < \infty$, $3 < q < \infty$, $0 < \alpha < \frac{1}{2}$
with $\frac{2}{s} + \frac{3}{q} = 1 - 2\alpha$ and $u_0 \in L_\sigma^2(\Omega) \cap \mathbb{B}_{q,s}^{-1+\frac{3}{q}}(\Omega)$.
Then the unique $L_\alpha^s(0, T; L^q)$ -strong solution u is *unique* on $[0, T)$
in the class of all well-chosen weak solutions.

Definition of well-chosen weak solutions I

Assumptions

- (i) Let $u_{0n} \rightharpoonup u_0$ in $L^2_\sigma(\Omega)$ and $u_{0n} \rightarrow u_0$ in $\mathbb{B}_{q,s}^{-1+\frac{3}{q}}(\Omega)$ as $n \rightarrow \infty$
- (ii) Let $(J_n) \subset \mathcal{L}(L^q_\sigma(\Omega), D(A_q^{1/2}))$ such that

$$\|J_n\|_{\mathcal{L}(L^q_\sigma)} + \left\| \frac{1}{n} A_q^{1/2} J_n \right\|_{\mathcal{L}(L^q_\sigma)} \leq C_q \text{ and } J_n u \xrightarrow{n \rightarrow \infty} u \text{ in } L^q_\sigma(\Omega)$$

Consider the **approximate (NSt)**

$$\begin{cases} \partial_t u_n - \Delta u_n + (J_n u_n) \cdot \nabla u_n + \nabla p_n = 0, & \operatorname{div} u_n = 0 \\ u_n|_{\partial\Omega} = 0, & u_n(0) = u_{0n} \end{cases}$$

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Example Yosida approximation $J_n = (I + \frac{1}{n} A^{1/2})^{-1}$, $u_{0n} = J_n u_0$

Consequences $\exists!$ weak solution $u_n \in \mathcal{LH}_\infty$:

$$\|u_n\|_{L^\infty(L^2)} + \|u_n\|_{L^2(H^1)} \leq C \|u_{0n}\|_2 \leq C < \infty$$

Definition of well-chosen weak solutions II

There exists $(u_{n_k}) \subset (u_n)$ and $v \in \mathcal{LH}_\infty$:

$$u_{n_k} \rightharpoonup v \text{ in } L^2(H_0^1), \quad u_{n_k} \overset{*}{\rightharpoonup} v \text{ in } L^\infty(L^2), \quad u_{n_k} \rightarrow v \text{ in } L^2(L^2).$$

$$\Rightarrow u_{n_k}(t_0) \rightarrow v(t_0) \text{ in } L^2(\Omega) \text{ for a.a. } t_0 \in (0, T).$$

$$\Rightarrow v \text{ is a weak solution of (NSt) with } v(0) = u_0$$

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$$\Rightarrow v \text{ is a weak solution of (NSt) with } v(0) = u_0$$

Definition The solution v constructed as above is called a *well-chosen weak solution* of (NSt) with $v(0) = u_0$.

Remark By the uniqueness theorem (to be proved) the well-chosen weak solution v with $v(0) = u_0$ is unique;
 $\Rightarrow v$ does *not* depend on the choice of $u_{0n} \rightarrow u_0$ and the choice of (weakly) convergent subsequences!

Proof of uniqueness theorem for well-chosen weak solutions I

(i) Construction of a local strong $L_\alpha^s(L^q)$ -solution u' on $[0, T']$:

$\exists \varepsilon_* > 0$, and since $u_{0n} \rightarrow u_0$ in $\mathbb{B}_{q,s}^{-1+3/q}(\Omega)$, we find $0 < T' \leq T$:

$$\int_0^{T'} (\tau^\alpha \|e^{-\tau A} u_{0n}\|_q)^s d\tau < \varepsilon_*$$

\Rightarrow the approximate (NSt) has a unique solution $u_n \in L_\alpha^s(0, T'; L^q)$ with bound

$$\|u_n\|_{L_\alpha^s(0, T'; L^q)} \leq C\varepsilon_* \quad \forall n$$

$\Rightarrow \exists$ subsequence (u'_{n_k}) of (u_{n_k}) : $u'_{n_k} \rightarrow u'$ in $L_\alpha^s(0, T'; L^q)$
 $\Rightarrow u' = v$ is a strong $L_\alpha^s(L^q)$ -solution on $[0, T']$ with $u'(0) = u_0$.

Proof of uniqueness theorem for well-chosen weak solutions II

Up to now we got 3 solutions:

$$u \in L_{\alpha}^s(0, T; L^q), \quad u' \in L_{\alpha}^s(0, T'; L^q), \quad v \in \mathcal{LH}_T$$

Uniqueness in $L_{\alpha}^s(0, T'; L^q) \Rightarrow u = u' = v$ on $(0, T')$

$\Rightarrow u_n \rightharpoonup u' = u$ in $L_{\alpha}^s(0, T'; L^q)$ (whole sequence (u_n)),

$u_n \rightharpoonup v = u$ in \mathcal{LH}_T

Proof of uniqueness theorem for well-chosen weak solutions II

Up to now we got 3 solutions:

$$u \in L_{\alpha}^s(0, T; L^q), \quad u' \in L_{\alpha}^s(0, T'; L^q), \quad v \in \mathcal{LH}_T$$

Uniqueness in $L_{\alpha}^s(0, T'; L^q) \Rightarrow u = u' = v$ on $(0, T')$

$\Rightarrow u_n \rightharpoonup u' = u$ in $L_{\alpha}^s(0, T'; L^q)$ (whole sequence (u_n)),

$u_n \rightharpoonup v = u$ in \mathcal{LH}_T

$\Rightarrow v$ satisfies (EE) on $[0, T']$ and, of course, (SEI) on $[0, T)$

$\Rightarrow v$ satisfies (EI) on $[T', T)$. Serrin's Uniqueness Theorem \Rightarrow

$$u = v \quad \text{on} \quad [T', T)$$



The limit $s = \infty$

Is the space $B_{q,\infty}^{-1+3/q}(\Omega)$ too large? For $\Omega = (0, 2\pi)^3_{\text{per}}$

$$\exists u_0 \in \bigcap_{1 < q \leq \infty} B_{q,\infty}^{-1+3/q}(\Omega) \cap L^2_\sigma(\Omega)$$

such that each Leray-Hopf weak solution u with $u(0) = u_0$ satisfies

$$\|u(t) - u_0\|_{B_{q,\infty}^{-1+3/q}} \geq \delta > 0$$

with δ independent of u , although

$$\|u(t) - u_0\|_{L^2(\Omega)} \rightarrow 0 \quad \text{as } t \rightarrow 0,$$

cf. Cheskidov-Shvydkoy, Proc. AMS 2010

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What about $B_{q,\infty}^{-1+3/q}(\Omega)^{\circ} := (\mathcal{D}(A_{q'})', L_{\sigma}^q(\Omega))_{1-\alpha,\infty}^{\circ}$?